

Time-stepping methods (continued) and numerical derivatives

(Pencil Code School, Geneva/CERN, 21st of September 2025)

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Overview:

- * Summary on common time-stepping methods
- * “Creative” methods for time stepping
- * Numerical derivatives
- * Numerical curiosities due to precision
- * Numerical curiosities from the initial condition
- * Numerical curiosities from “switching on” the simulation

Time integration methods

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Euler method:

- * very simple to implement (single-step integration)
- * may have infinitely growing error

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- * very accurate and computation efficient (larger Δt is possible)

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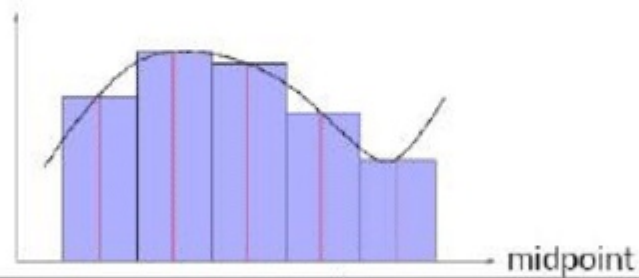
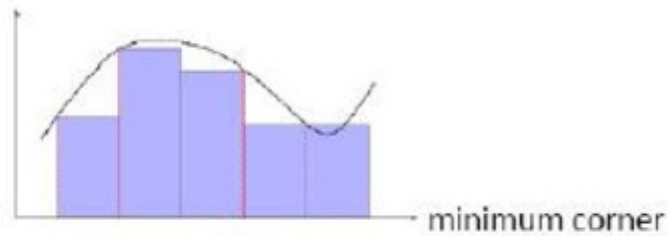
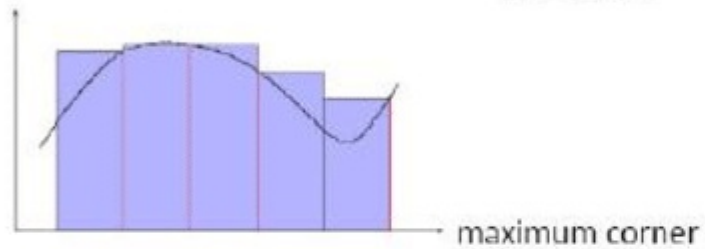
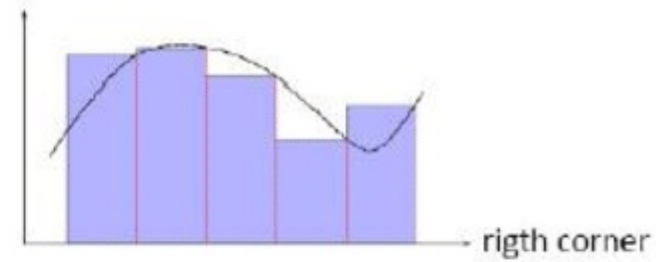
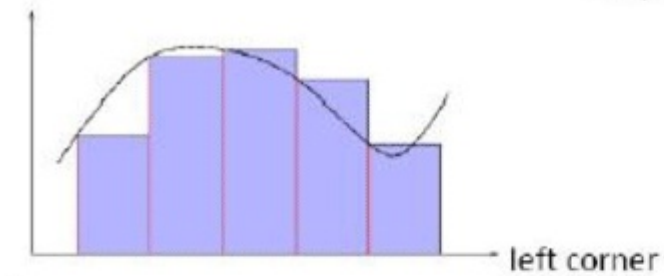
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Runge-Kutta scheme:

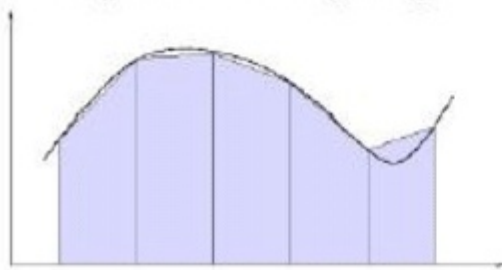
- * harder to implement (order of accuracy = number of integration substeps)
- * very accurate and computation efficient (larger Δt is possible)
- * substeps require additional function evaluations and derivatives (\Rightarrow costs)

Time integration methods

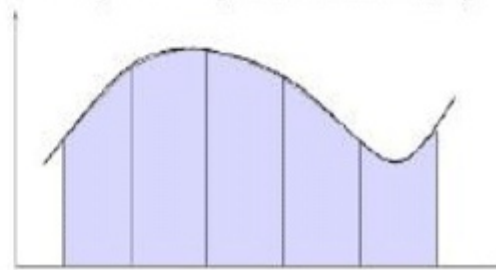
Rectangle
Step width * height



Trapezoid
Step width * average height

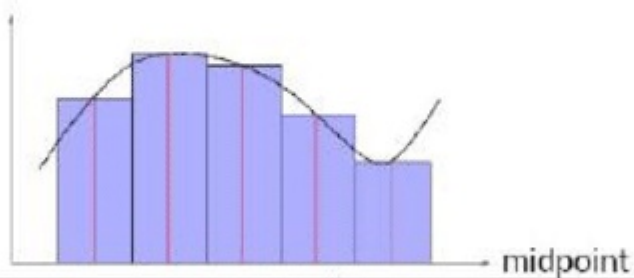
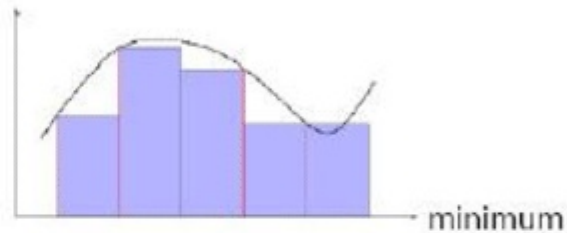
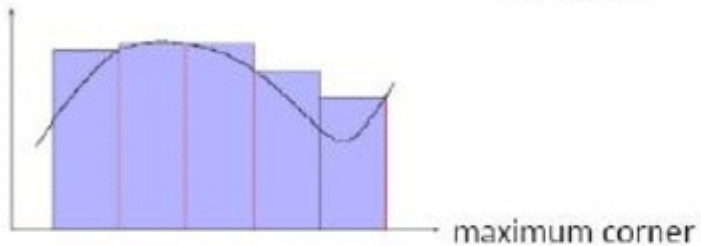
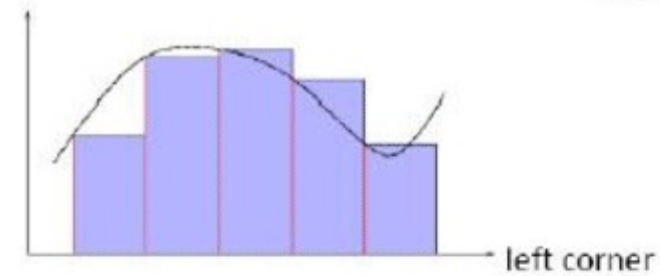


Simpson's
Step width * Simpson's height

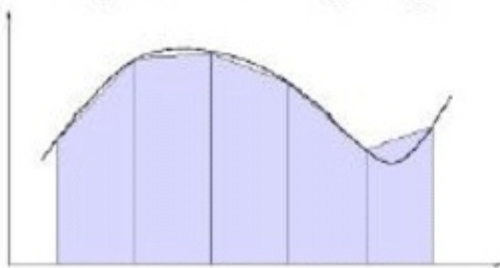


Time integration methods

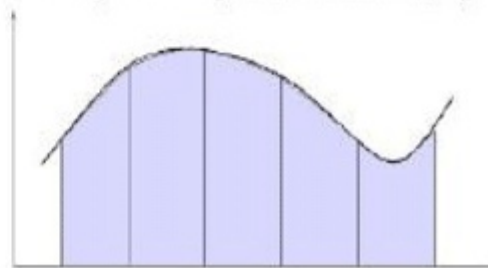
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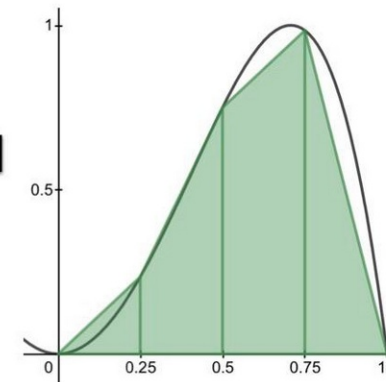
Trapezoid
Step width * average height



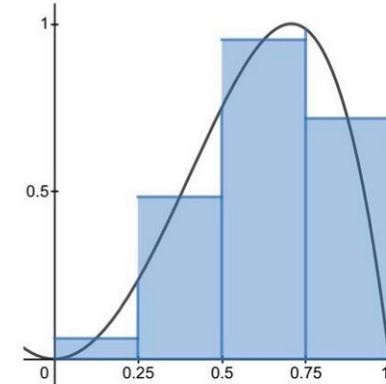
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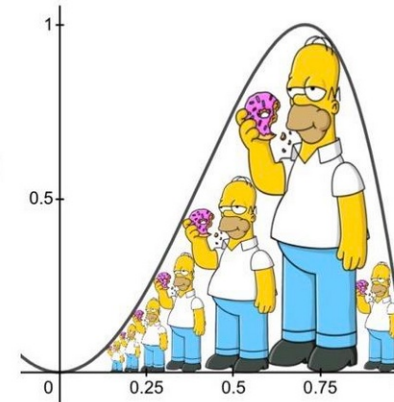
Trapezoidal Rule



Midpoint Rule



Simpson's Rule



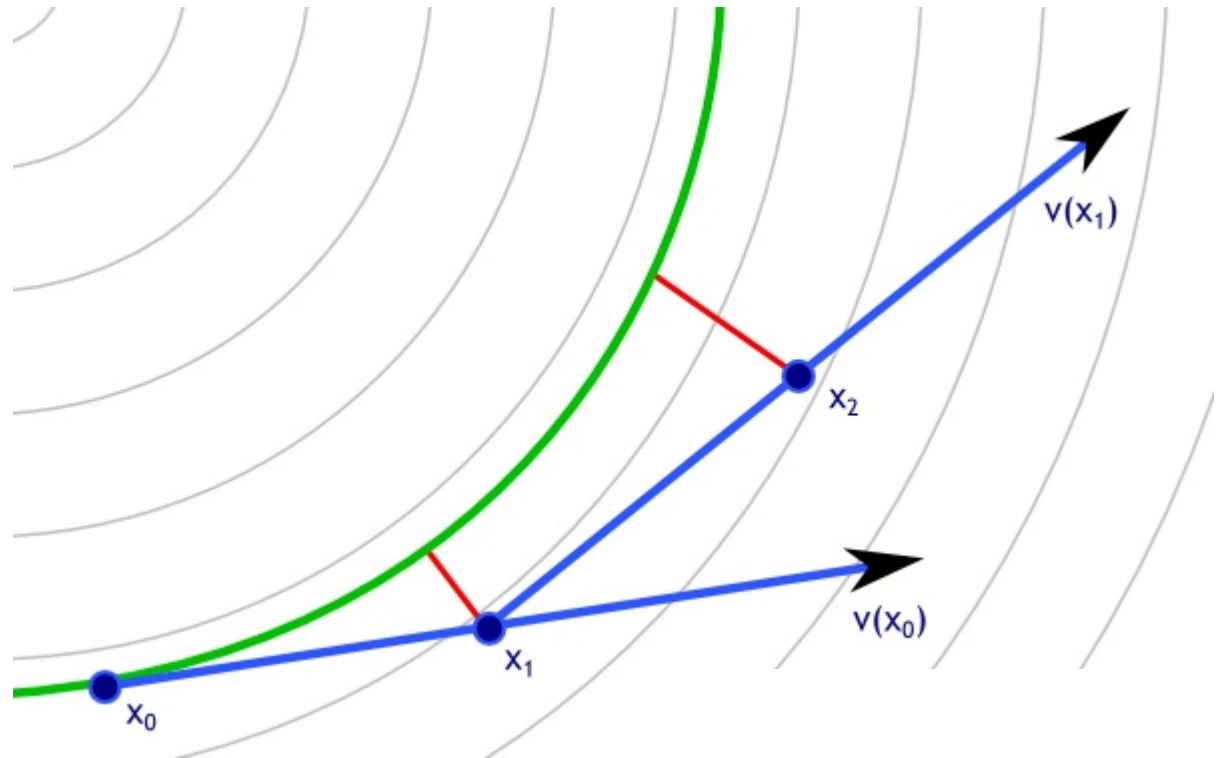
Euler method

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- 1) start from x_0
- 2) calculate derivatives
(tangential)

gray: true field

green: exact solution



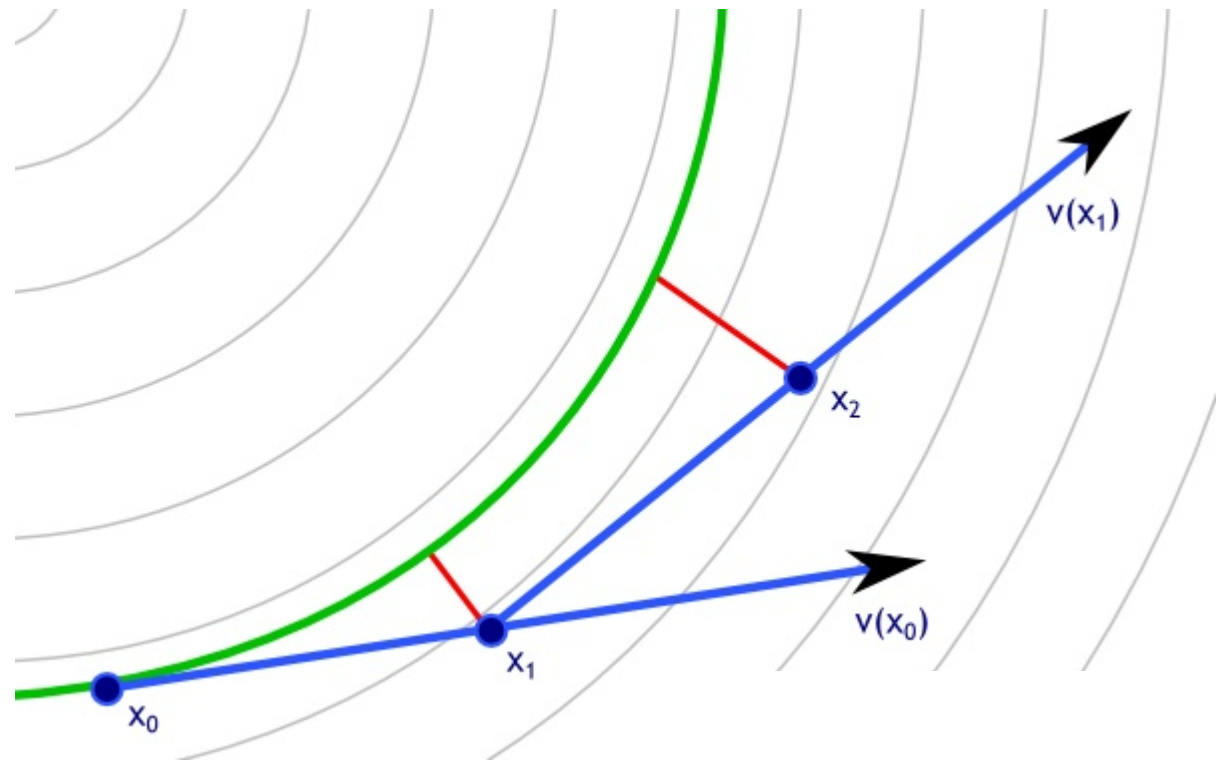
Euler method

- 1) start from x_0
- 2) calculate derivatives
(tangential)
- 3) iterate to x_1
with $dt=0.5$

green: exact solution

blue: tangential derivative

red: deviation



$$x_{n+1} = x_n + dt \cdot v(x_n) + O(dt^2)$$

x_n : current position

x_{n+1} : next position

dt : time step

$v(x_n)$: vector field at position x_n

$O(dt^2)$: error

=> error grows infinitely...

=> make dt smaller?

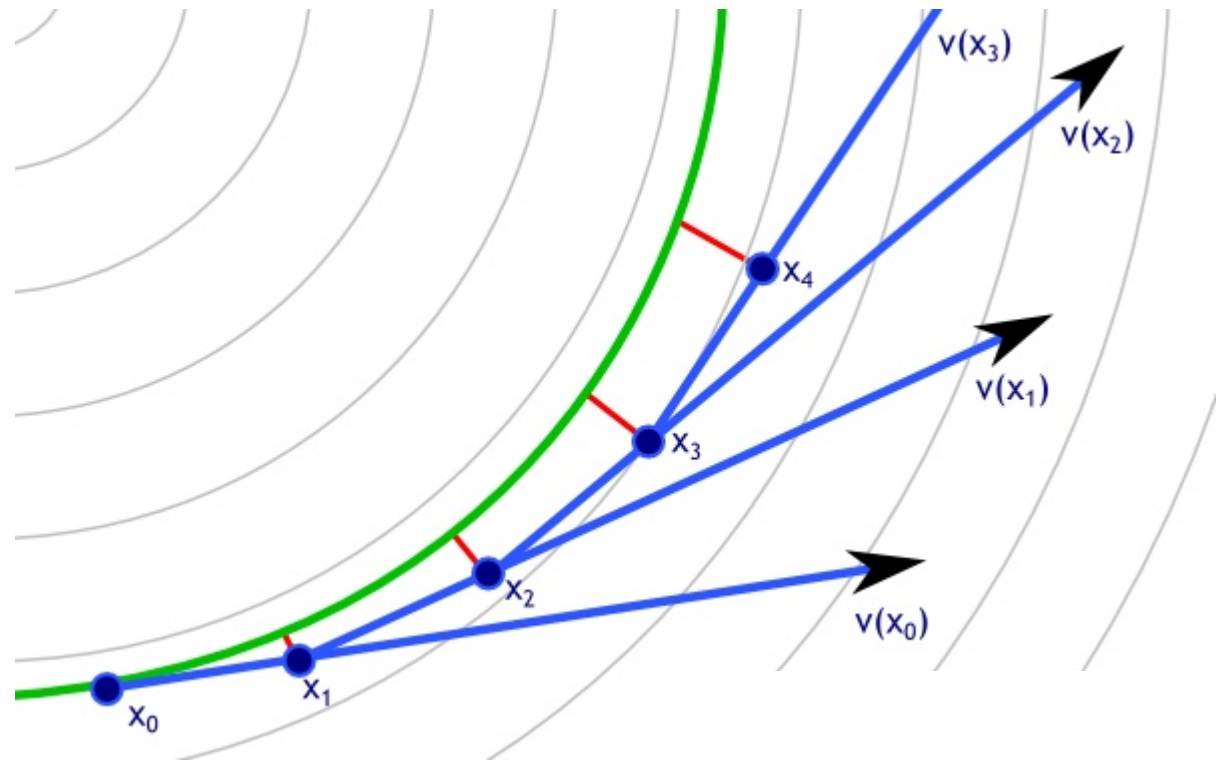
Euler method

- 1) start from x_0
- 2) calculate derivatives
(tangential)
- 3) iterate to x_1
with $dt=0.25$

green: exact solution

blue: tangential derivative

red: deviation



$$x_{n+1} = x_n + dt \cdot v(x_n) + O(dt^2)$$

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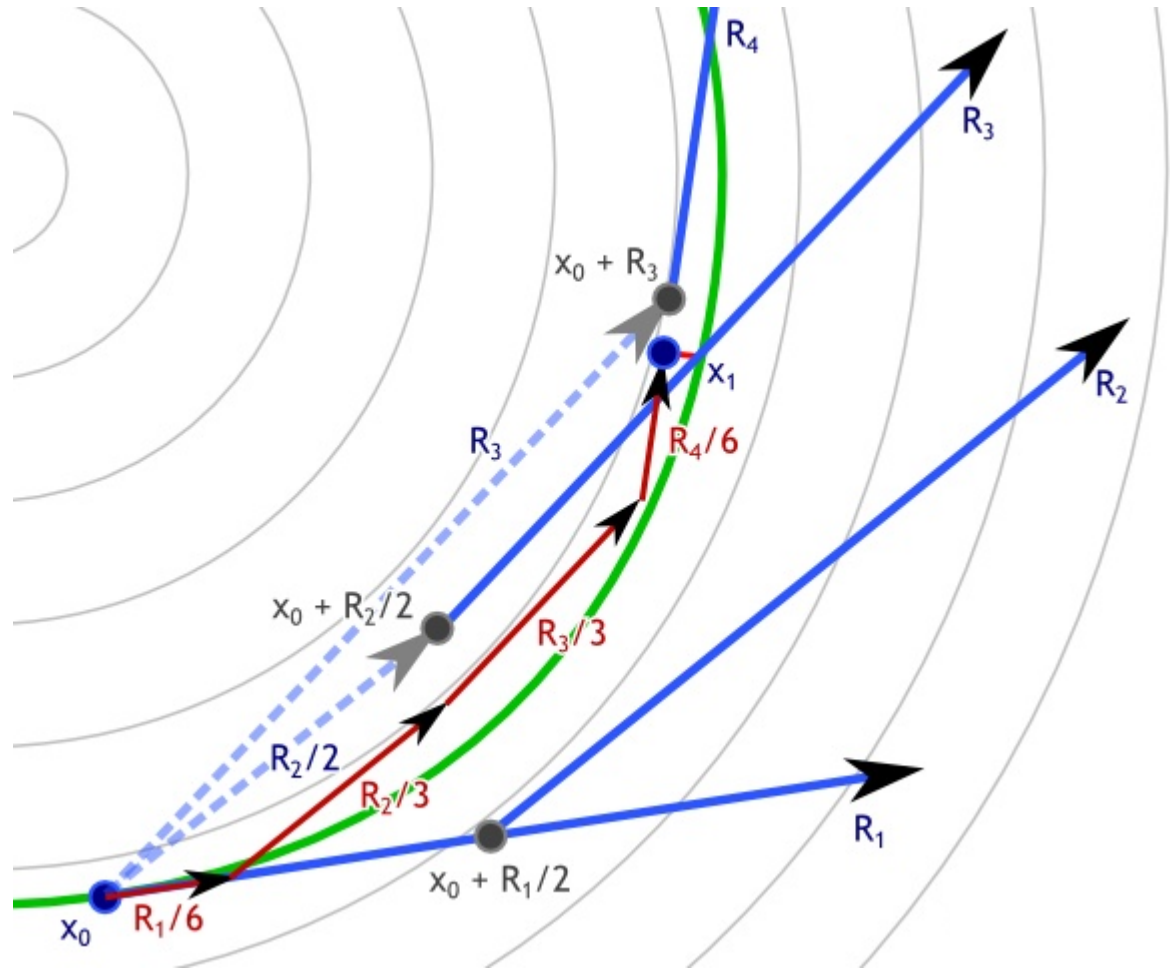
$O(dt^2)$: error

=> accuracy can be improved
by making dt smaller
=> but error still grows...

Runge-Kutta method (4th order)

- 1) start from x_0
- 2) calculate derivatives (tangential)
- 3) iterate substeps and recalculate derivatives (tangential)
- 4) reach x_1 with substeps

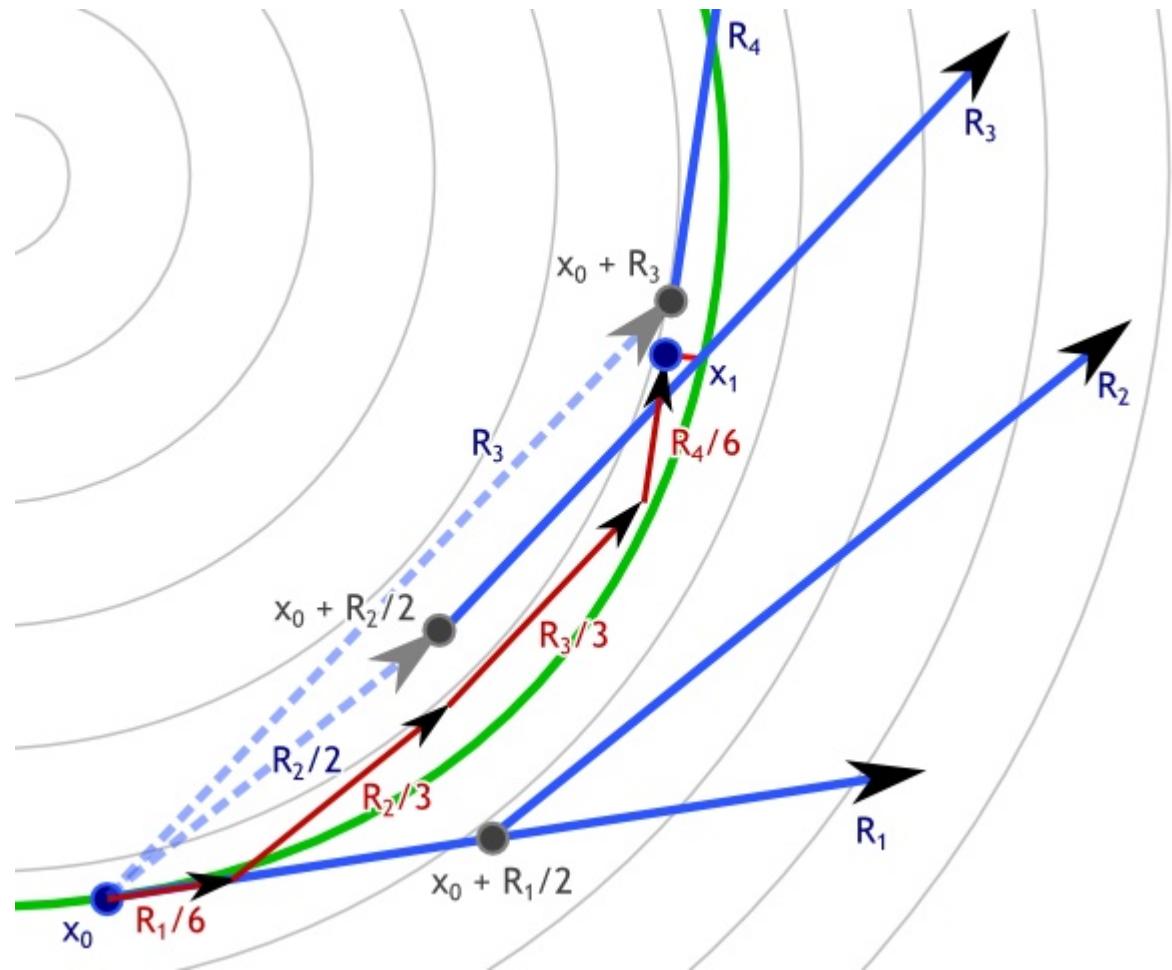
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blue: tangential derivative
red: deviation



=> much better accuracy!

=> larger computational cost

Runge-Kutta method (4th order)

RK4 scheme:

$$k_1 = dt \cdot v(x_n)$$

$$k_2 = dt \cdot v\left(x_n + \frac{k_1}{2}\right)$$

$$k_3 = dt \cdot v\left(x_n + \frac{k_2}{2}\right)$$

$$k_4 = dt \cdot v(x_n + k_3)$$

$$x_{n+1} = x_n + \frac{k_1}{6} + \frac{k_2}{3} + \frac{k_3}{3} + \frac{k_4}{6} + O(dt^5)$$

x_n : current position

x_{n+1} : next position

dt : time step

$v(x)$: vector field at position x

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=> effect: back-reaction to system is delayed => slower overall dynamics!

more timesteps needed & finally not such a big saving of computing time.

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=> Speed up and accuracy granted only in 1D.

“Creative” time-stepping methods

⇒ Split Runge-Kutte time step into two halves (**timestep_strang**).

Idea: ...?

=> effects?

“Creative” time-stepping methods

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Idea: ...?

=> effects?

➡ Algorithms for stiff PDEs (**timestep_stiff** from Numerical Recipies).

Idea: ...?

=> effects?

Numerical derivatives (6th order)

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Pencil Code manual:

H Numerical methods

H.1 Sixth-order spatial derivatives

Spectral methods are commonly used in almost all studies of ordinary (usually incompressible) turbulence. The use of this method is justified mainly by the high numerical accuracy of spectral schemes. Alternatively, one may use high order finite differences that are faster to compute and that can possess almost spectral accuracy. Nordlund & Stein [32] and Brandenburg et al. [16] use high order finite difference methods, for example fourth and sixth order compact schemes [28].¹⁹

The sixth order first and second derivative schemes are given by

$$f'_i = (-f_{i-3} + 9f_{i-2} - 45f_{i-1} + 45f_{i+1} - 9f_{i+2} + f_{i+3})/(60\delta x), \quad (235)$$

$$f''_i = (2f_{i-3} - 27f_{i-2} + 270f_{i-1} - 490f_i + 270f_{i+1} - 27f_{i+2} + 2f_{i+3})/(180\delta x^2), \quad (236)$$

Numerical derivatives (6th order)

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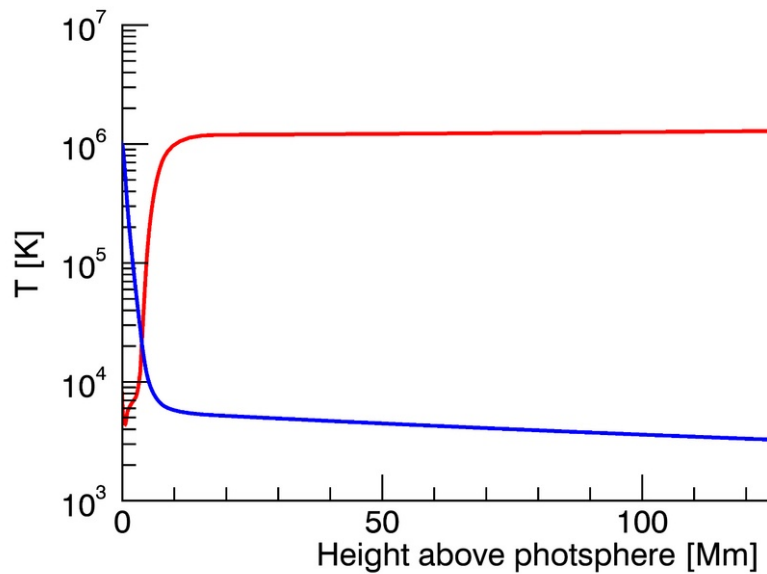
=> Coefficients can be obtained analytically.

Resolving steep gradients in the solar corona

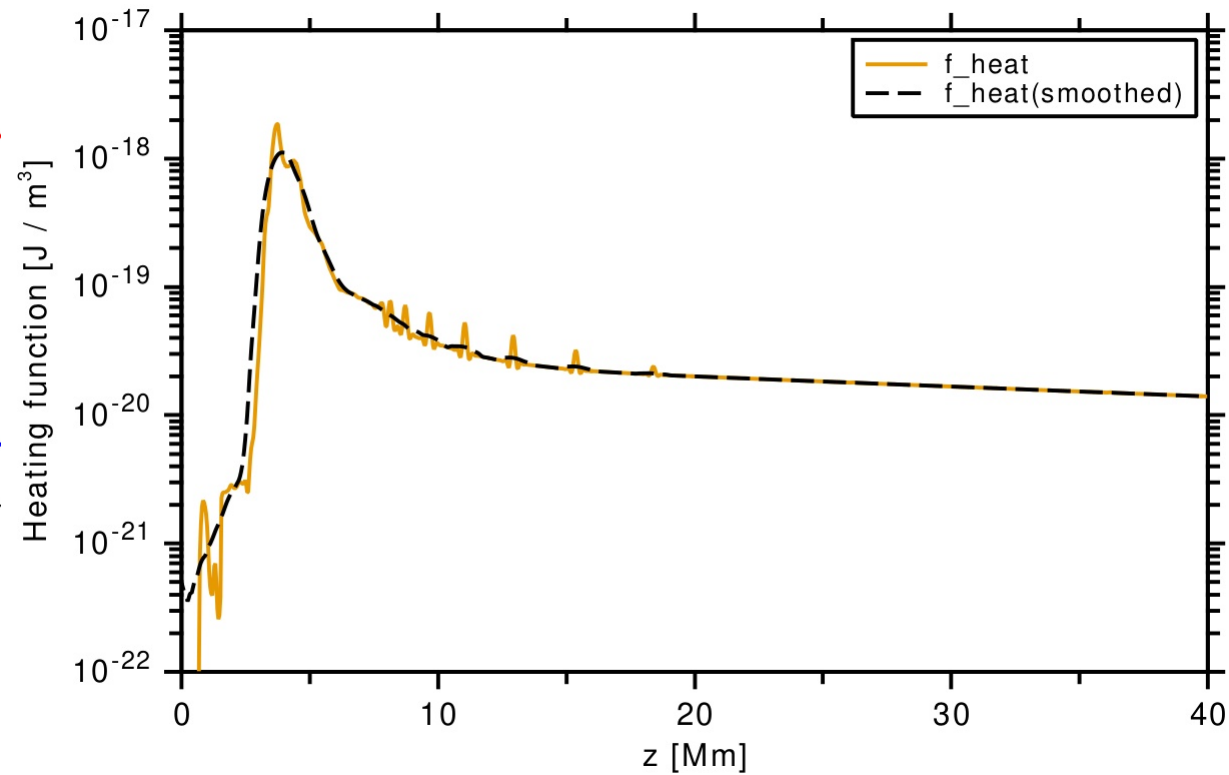
(Vartika Pandey)

Coronal heating in 1D MHD models:

- Initial condition:

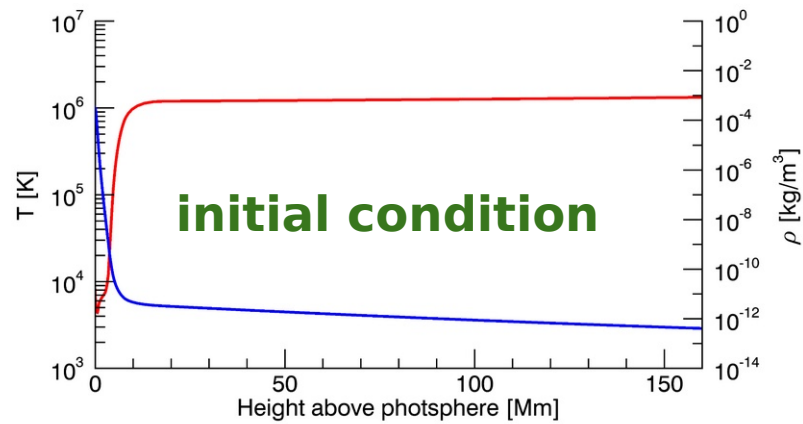


=> After **one** iteration:

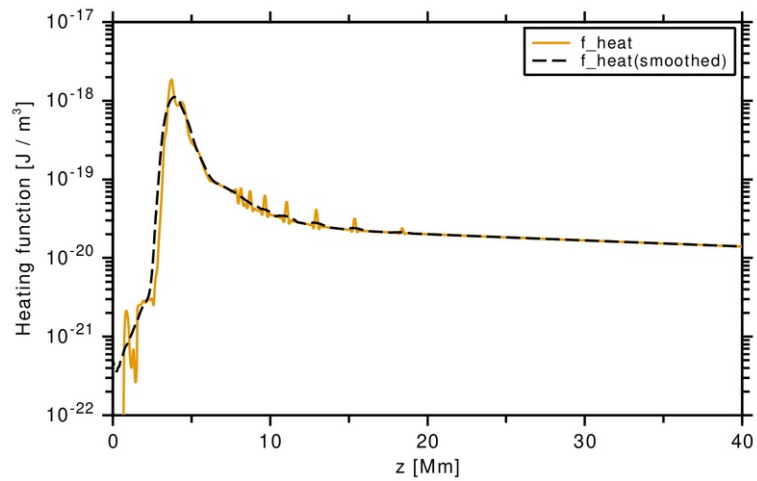


(Pandey & Bourdin, 2025)

Numerical stability analysis:

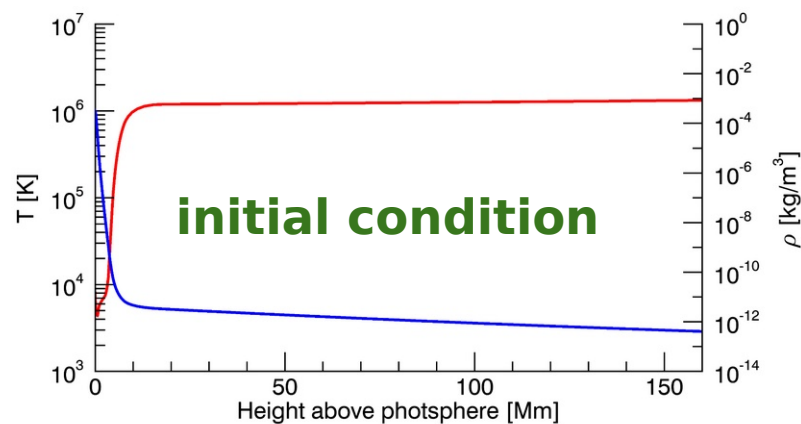


+ heating

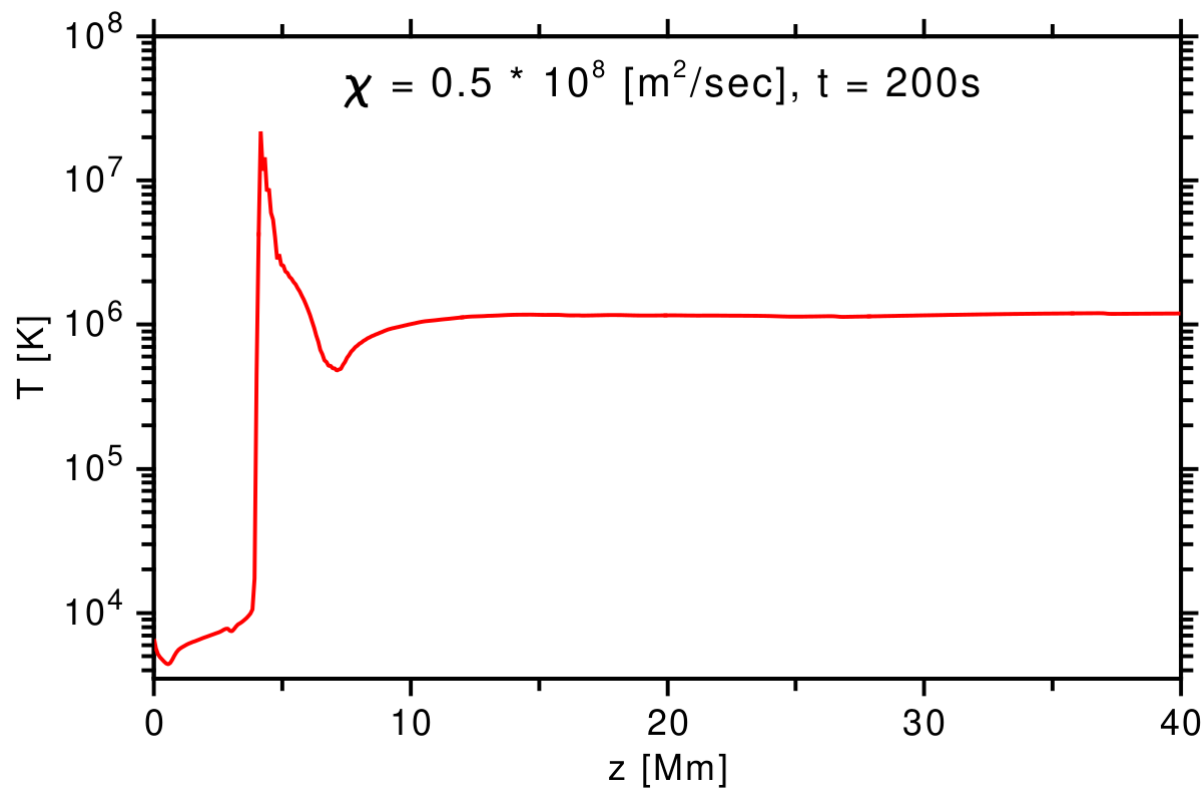
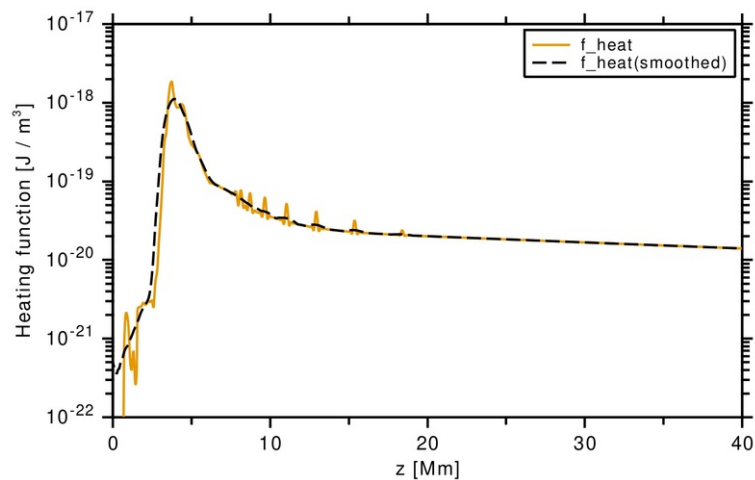


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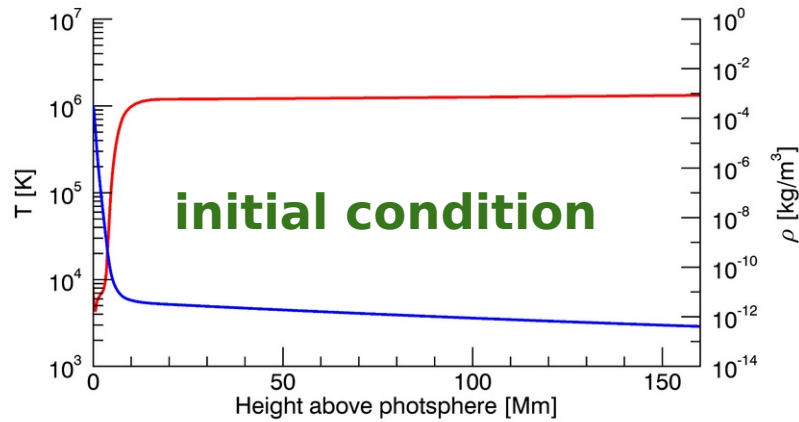


+ heating \Rightarrow

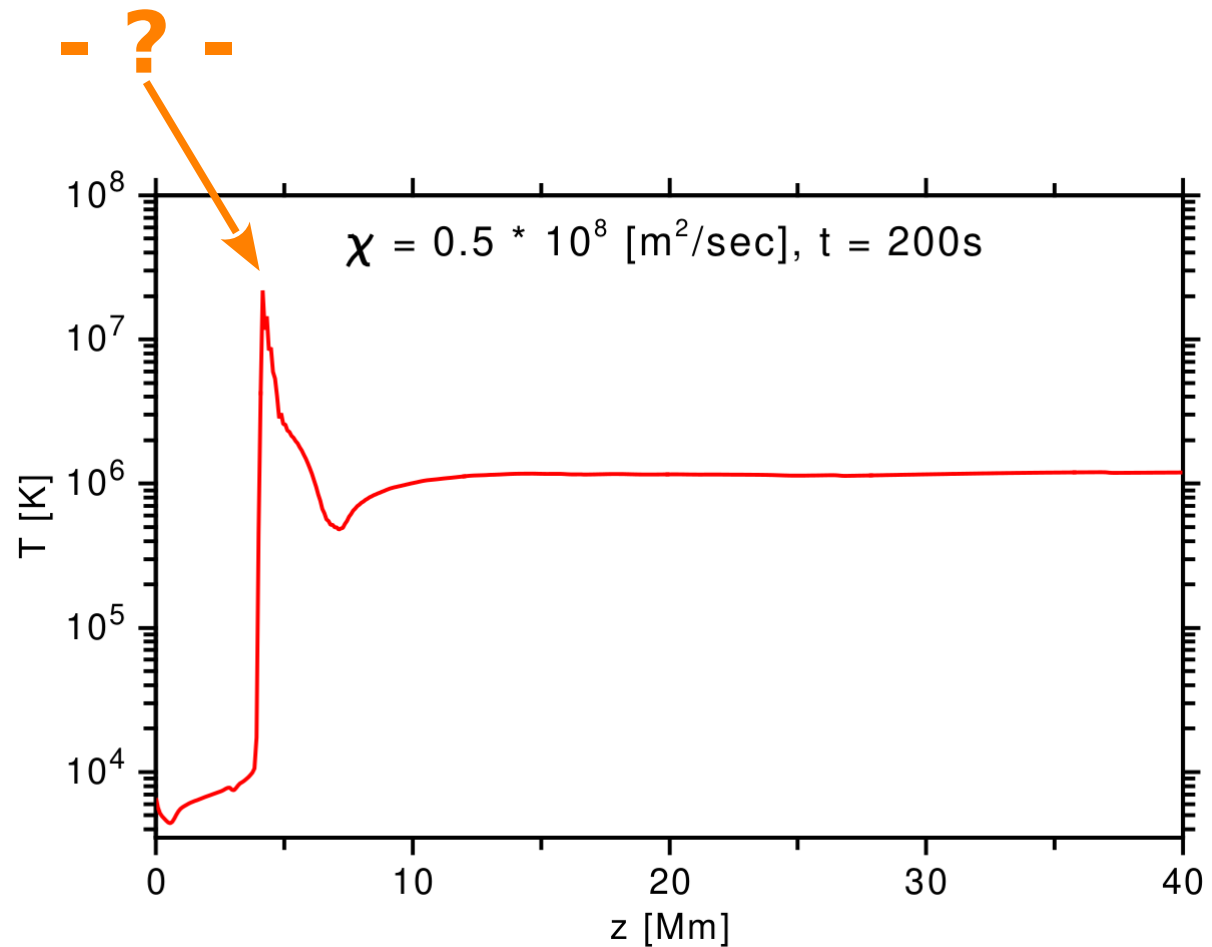
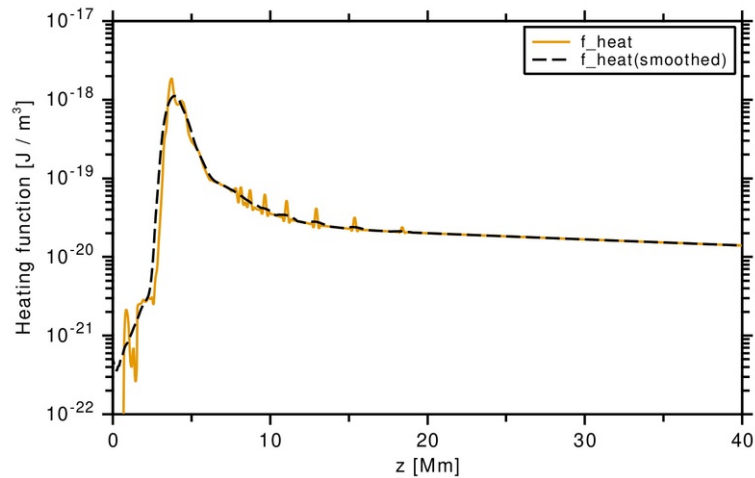


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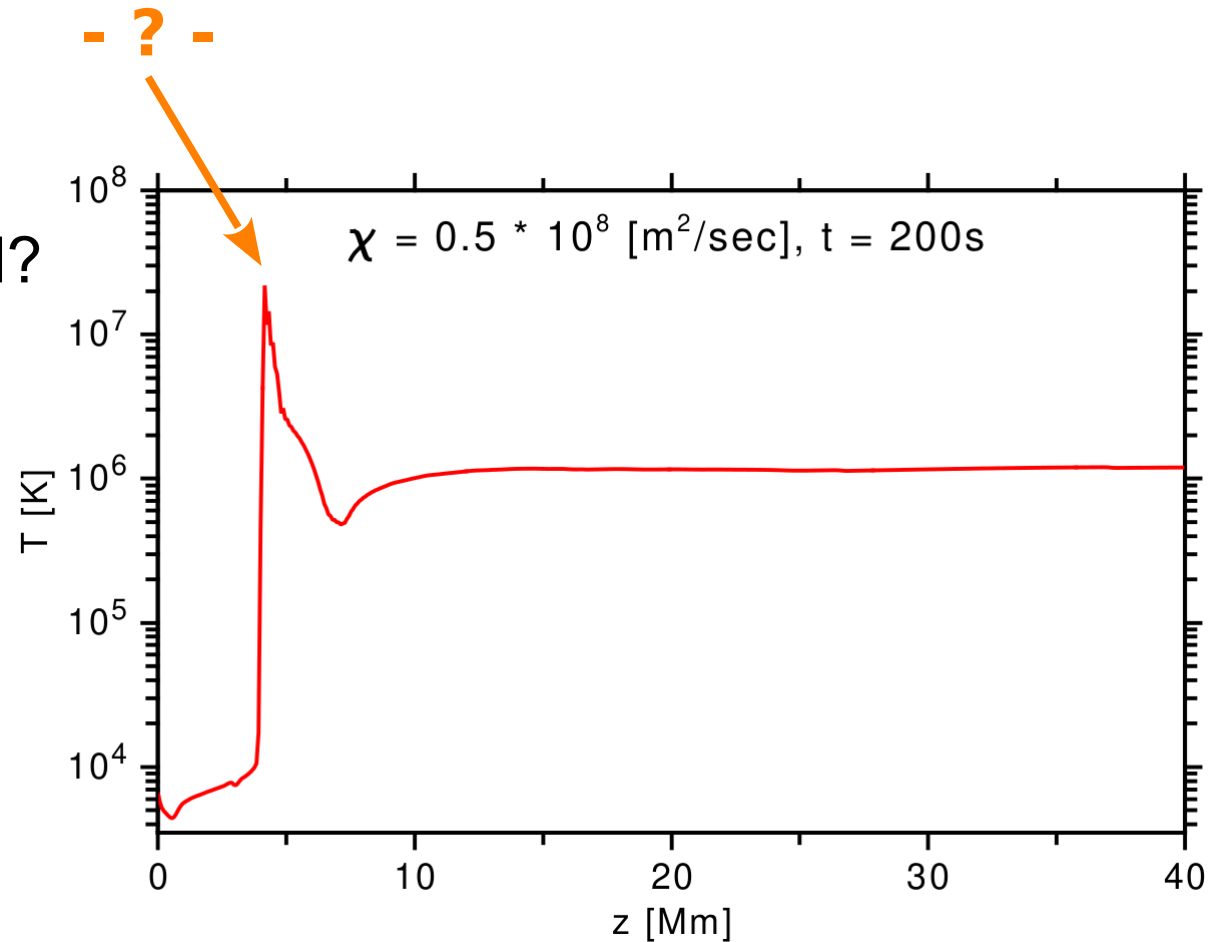


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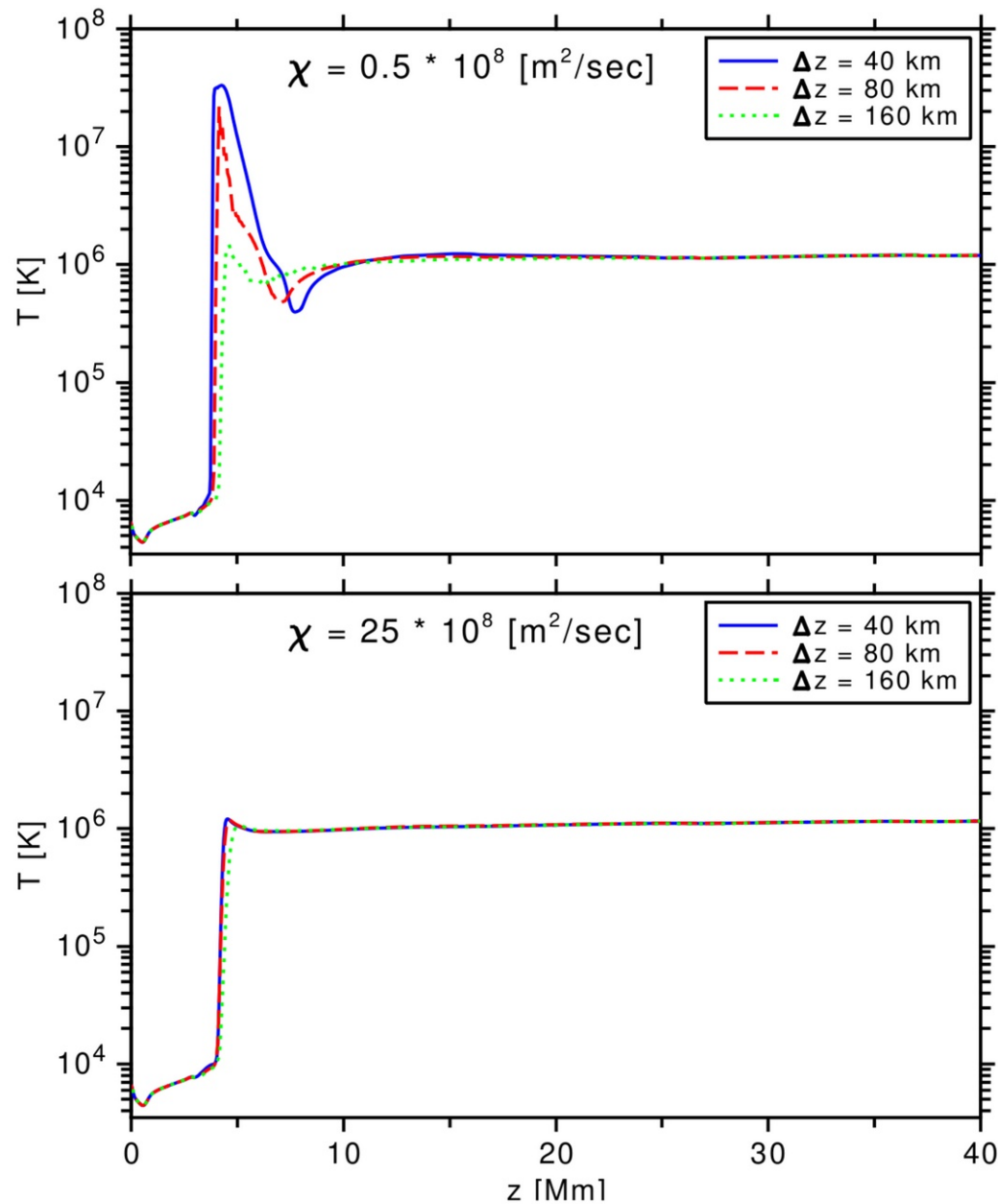
Possible problems:

- 1) too coarse grid?
- 2) insufficient diffusion?
- 3) absence of magnetic field?



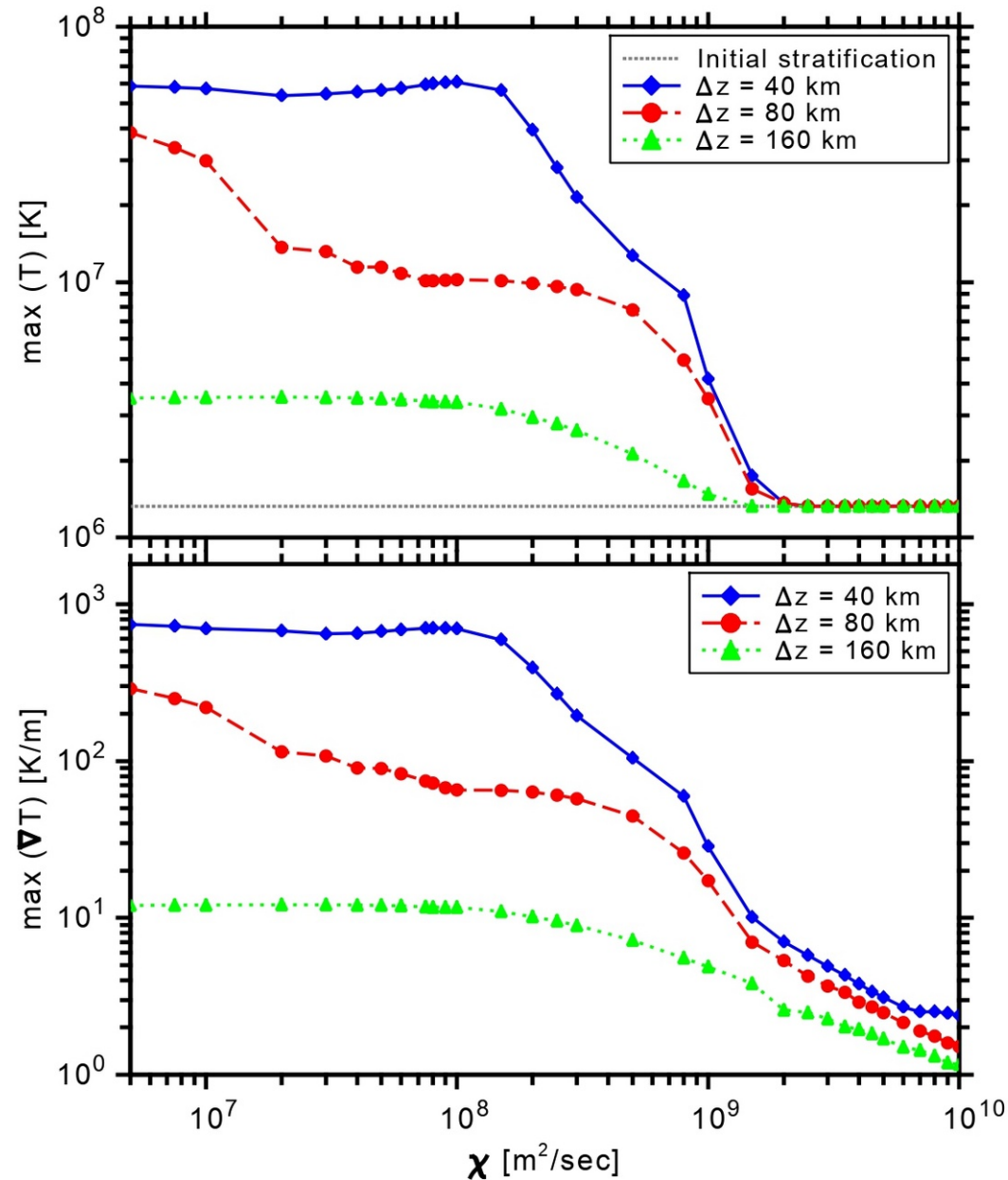
(Pandey & Bourdin, 2025)

Effects due to the grid resolution and heat conduction:



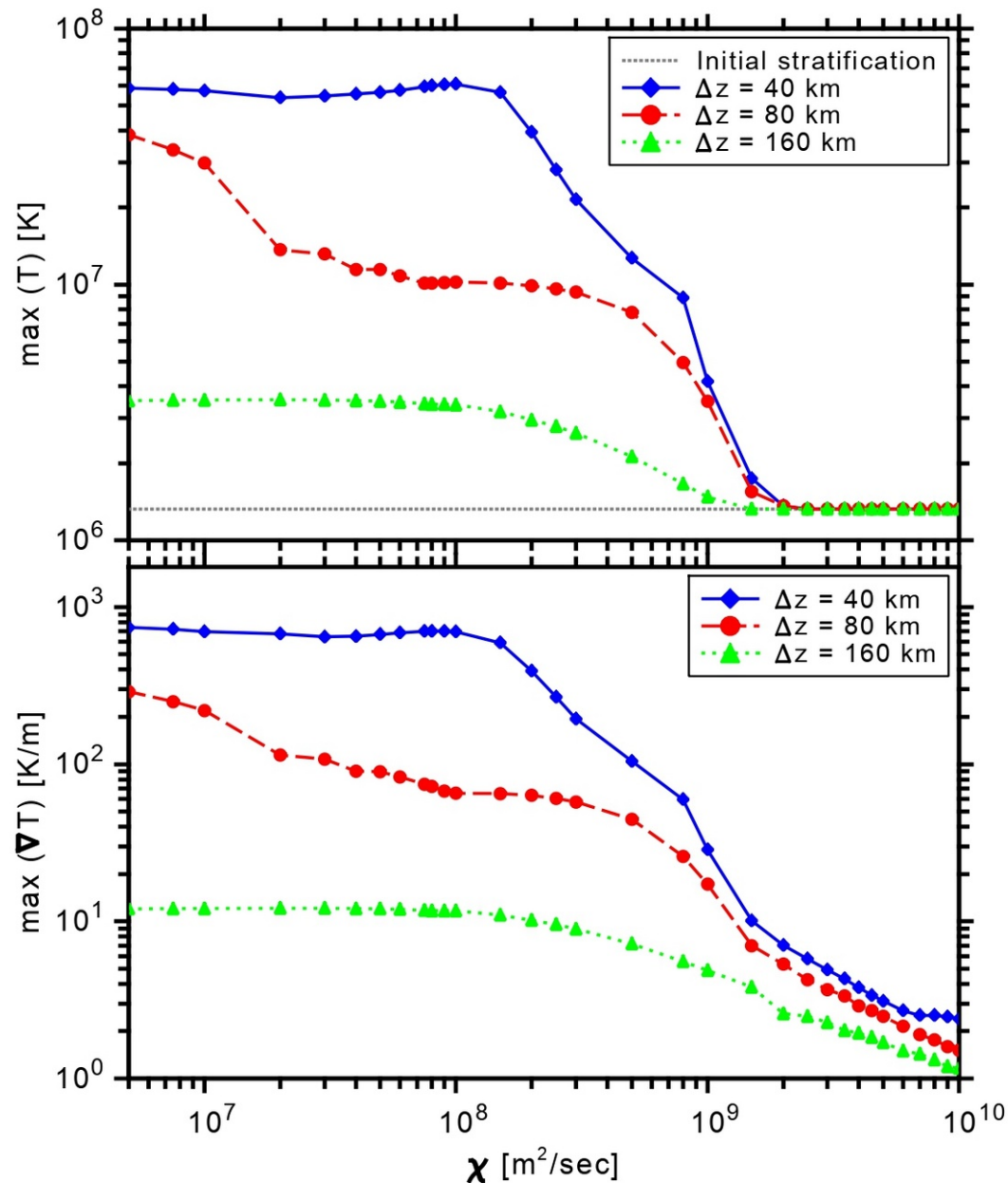
(Pandey & Bourdin, 2025)

Scaling the heat conduction:



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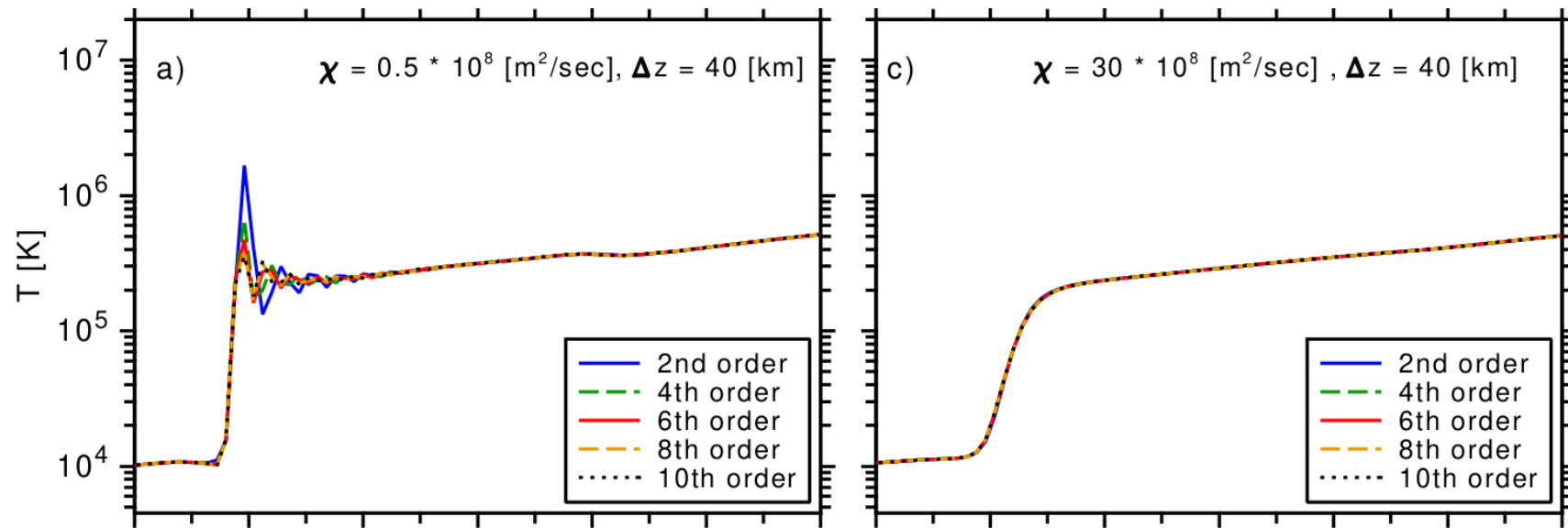


=> higher resolutions
do not always imply
models are “better”

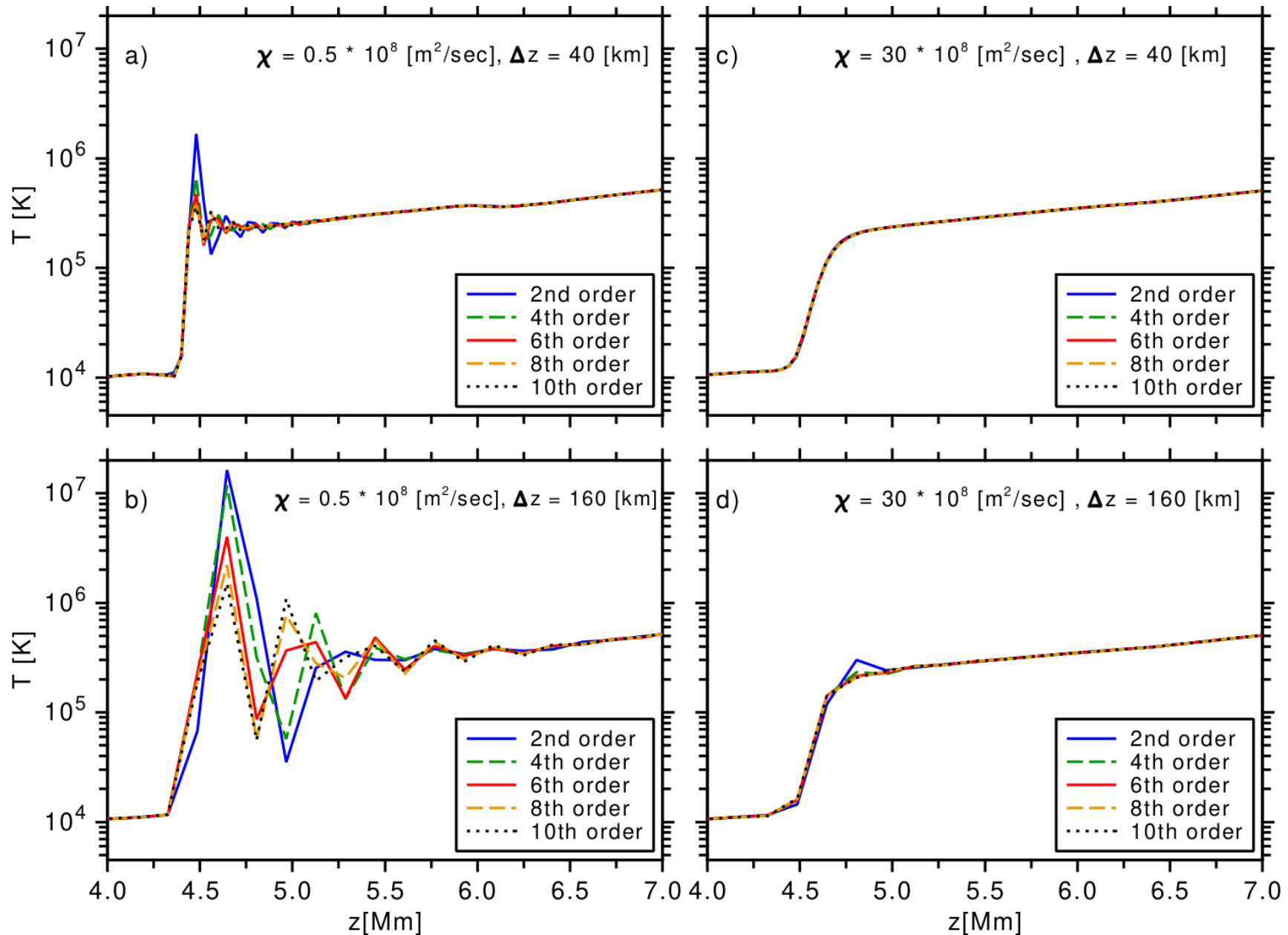
=> diffusion constants
need to fit to the
grid resolution

(Pandey & Bourdin, 2025)

Effects due to the order of numerical derivatives:



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Numerical curiosities (due to precision & truncation errors)

Precision and machine epsilon

Numerical curiosities cabinet:

$$(x + y) + z \neq x + (y + z)$$

Precision and machine epsilon

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Machine epsilon:

$$1.0 + \epsilon = 1.0$$

Precision and machine epsilon

Numerical curiosities cabinet:

$$(x + y) + z \neq x + (y + z)$$

Machine epsilon:

$$1.0 + \epsilon = 1.0 \quad (*)$$

- Estimation of the granularity:

($b = 2$; $p_0 = 0.5$)

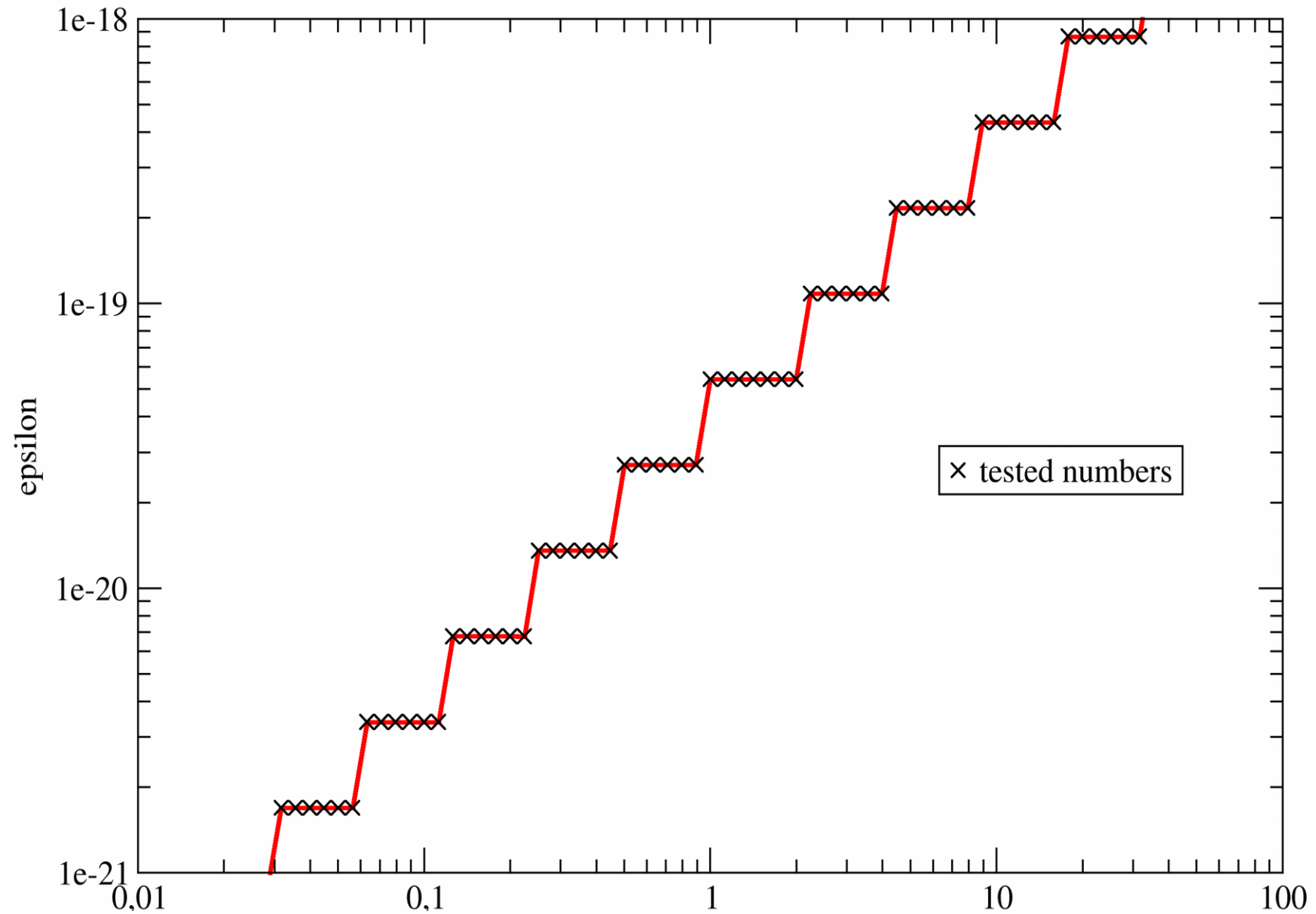
$$\epsilon_N = \frac{b}{2} \cdot b^{-p_N}$$

$$p_{N+1} = p_N \cdot b$$

Stop, if (*) is true

Precision and machine epsilon

Machine epsilon in Fortran 90 (double precision):



Granularity of numbers

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Numerical curiosities cabinet: (continued...)

Example: simple increment of simulation time:

$$T_N + dT = T_{N+1}$$

Granularity of numbers

Numerical curiosities cabinet: (continued...)

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$$4096 + 4 \cdot 10^{-4} = 4096$$

Granularity of numbers

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Example: simple increment of simulation time:

$$T_N + dT = T_{N+1}$$

$$4096 + 4 \cdot 10^{-4} = 4096$$

$$\Rightarrow \frac{T_N}{dT} \approx 10^{-7}$$

Granularity of numbers

Numerical curiosities cabinet: (continued...)

Example: simple increment of simulation time:

$$T_N + dT = T_{N+1}$$

$$4096 + 4 \cdot 10^{-4} = 4096$$

$$\Rightarrow \frac{T_N}{dT} \approx 10^{-7}$$

But:

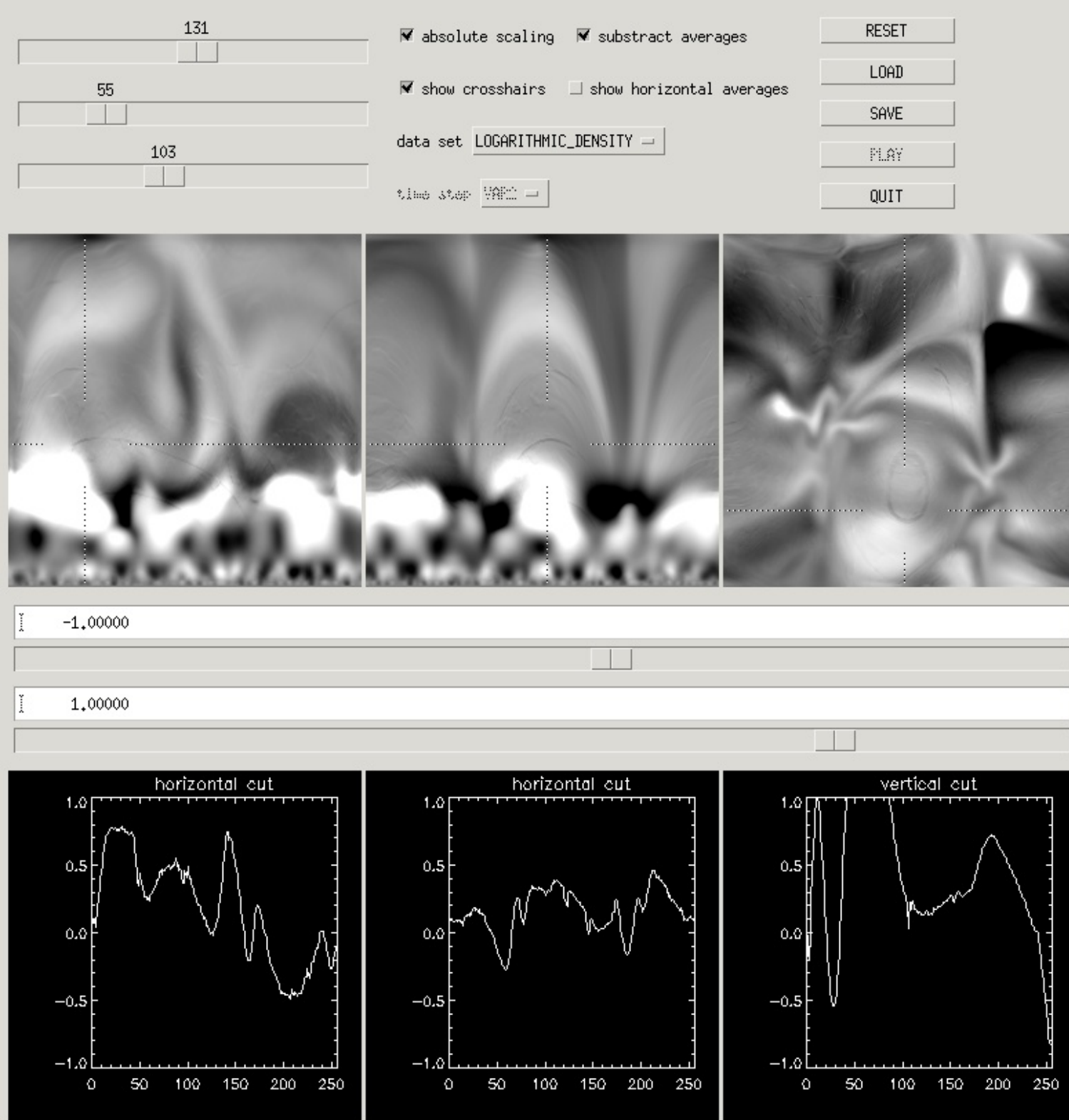
$$4096 + 8 \cdot 10^{-4} \neq 4096$$

Granularity effects in physical observables

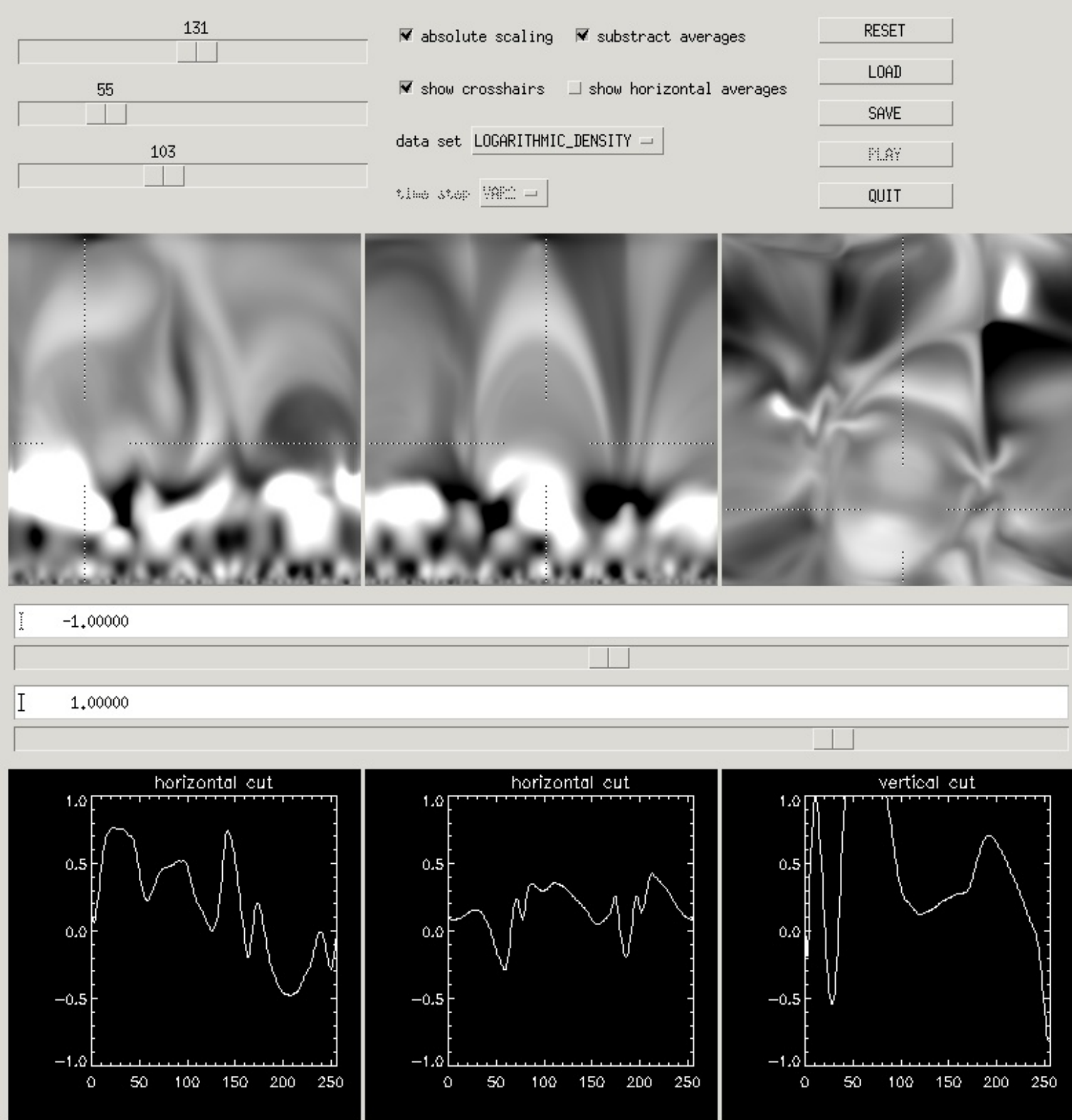
Granularity effects in physical observables

Density in a coronal 3D MHD simulation:

In(p):
SP



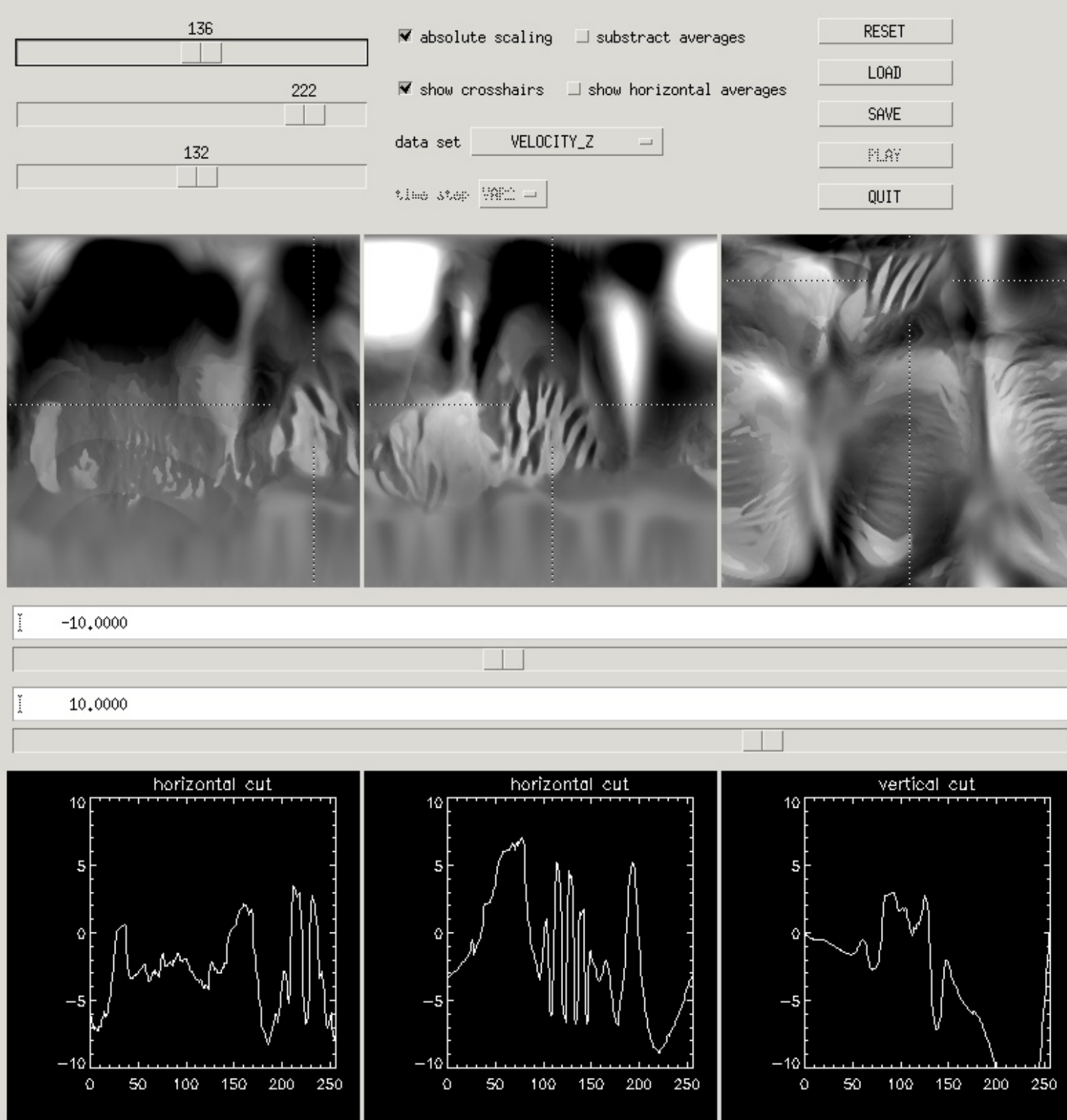
In(p):
DP



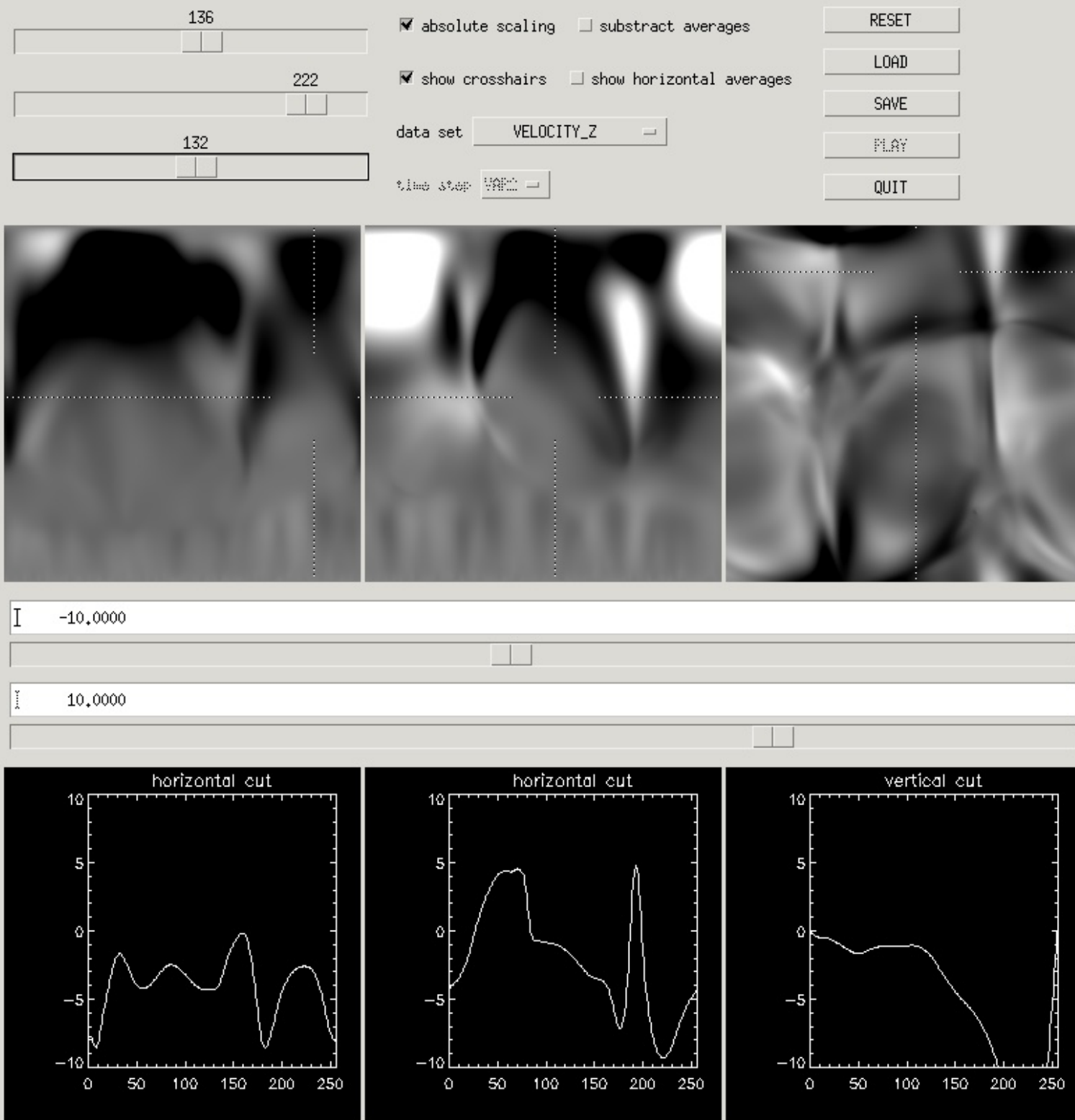
Granularity effects in physical observables

Vertical velocity:

Vz:
SP



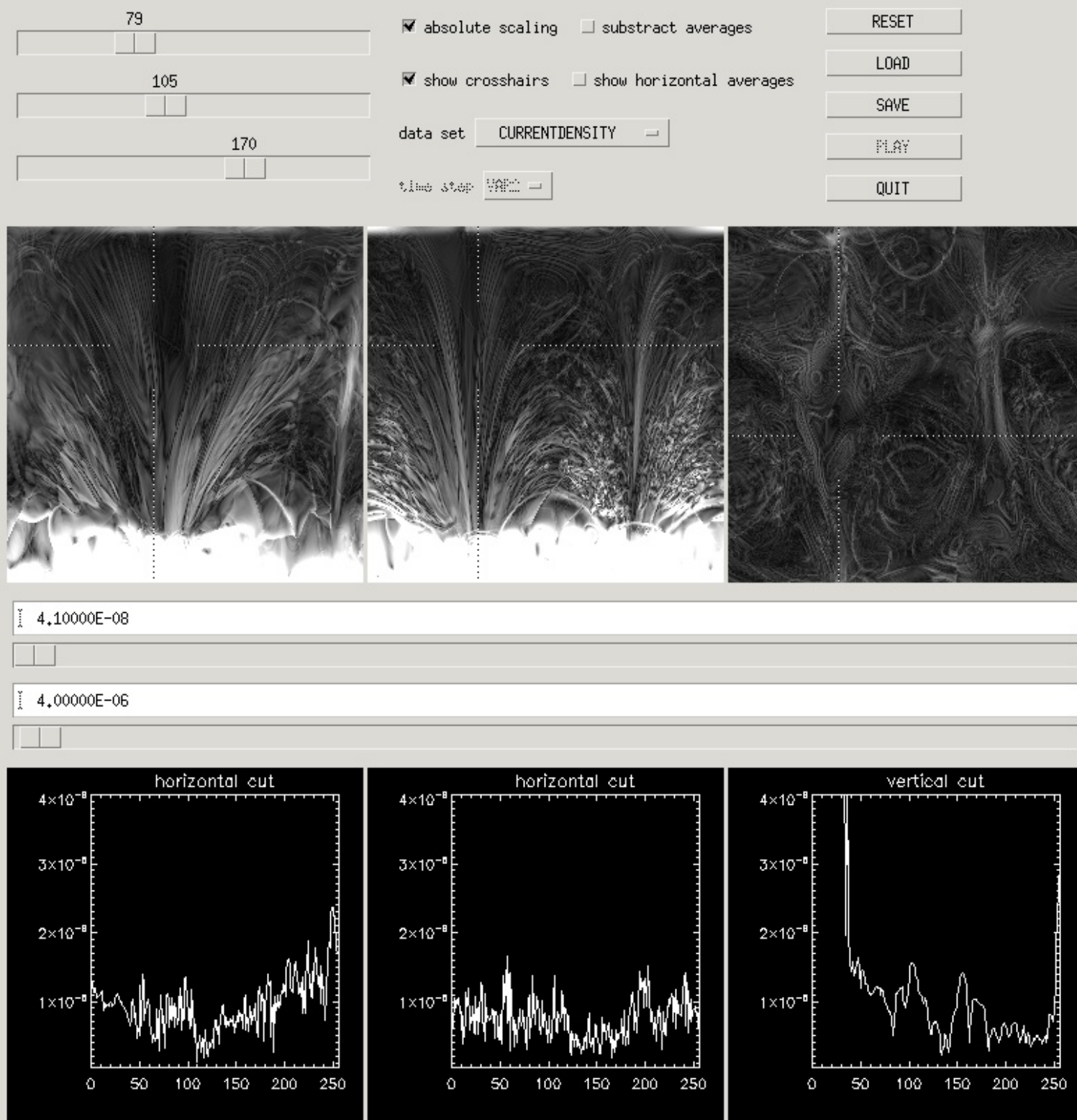
Vz:
DP



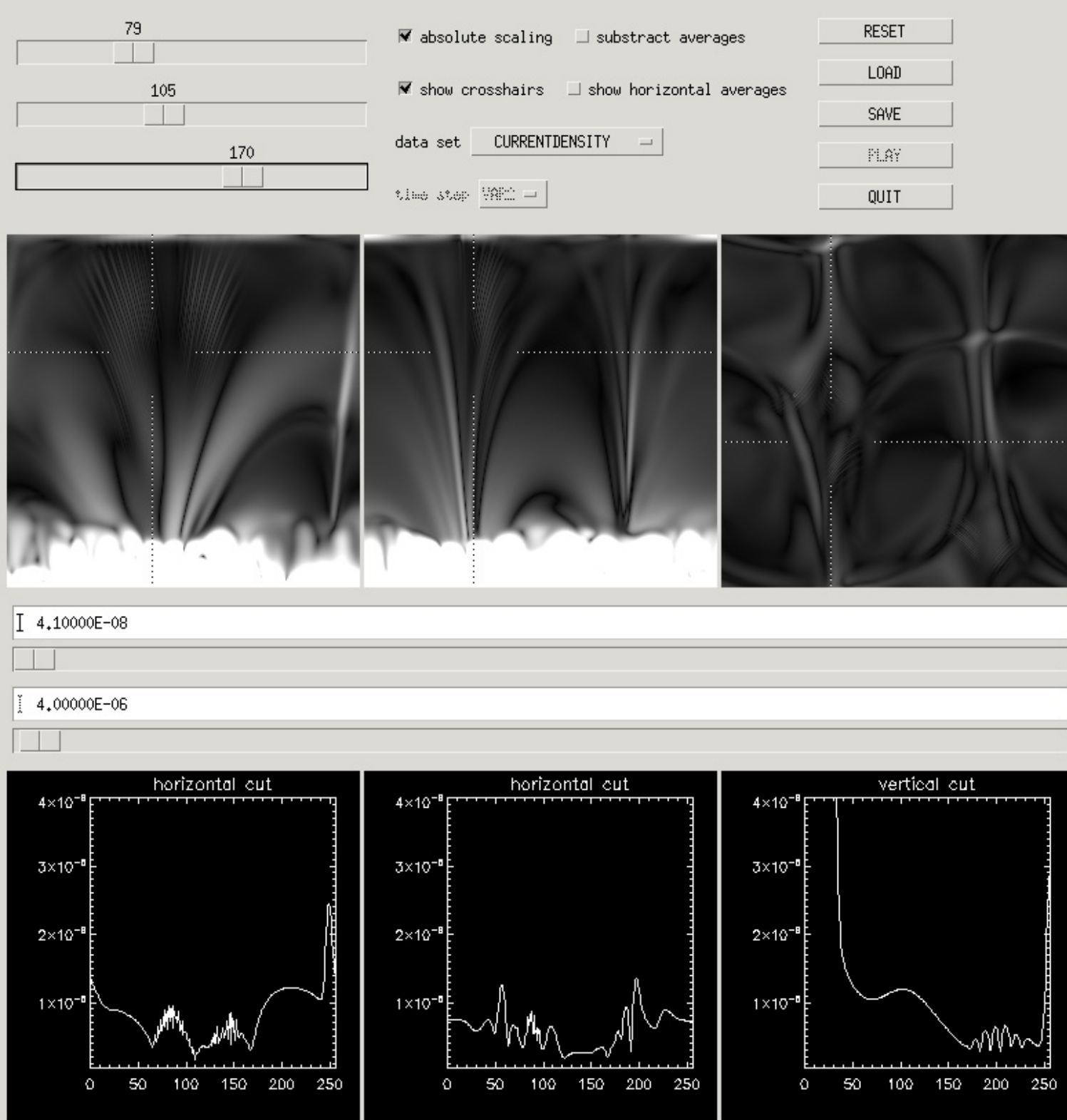
Granularity effects in physical observables

Current density / Heating rate: $H = \mu_0 \eta \mathbf{j}^2 \sim \nabla \times (\nabla \times A)$

|j|:
SP



|j|:
DP



Granularity effects in physical observables

Catastrophic cancellation:

“What Every Computer Scientist Should Know About Floating-Point Arithmetic”

(David Goldberg, 1991)

$$\nabla \times A \sim B \quad ; \quad \nabla \times B \sim j$$

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$$j_z \sim \frac{\partial B_y}{\partial x} - \frac{\partial B_x}{\partial y}$$

$$B_x \approx B_y$$

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$$B_x \approx B_y \quad \Rightarrow \quad \partial B_x \approx \partial B_y$$

Granularity effects in physical observables

Catastrophic cancellation:


“What Every Computer Scientist Should Know About Floating-Point Arithmetic”

(David Goldberg, 1991)

$$\nabla \times A \sim B \quad ; \quad \nabla \times B \sim j$$

$$j_z \sim \frac{\partial B_Y}{\partial x} - \frac{\partial B_X}{\partial y}$$

$$B_X \approx B_Y \quad \Rightarrow \quad \partial B_X \approx \partial B_Y$$

$$\Rightarrow j_z \approx 2 \cdot \epsilon$$


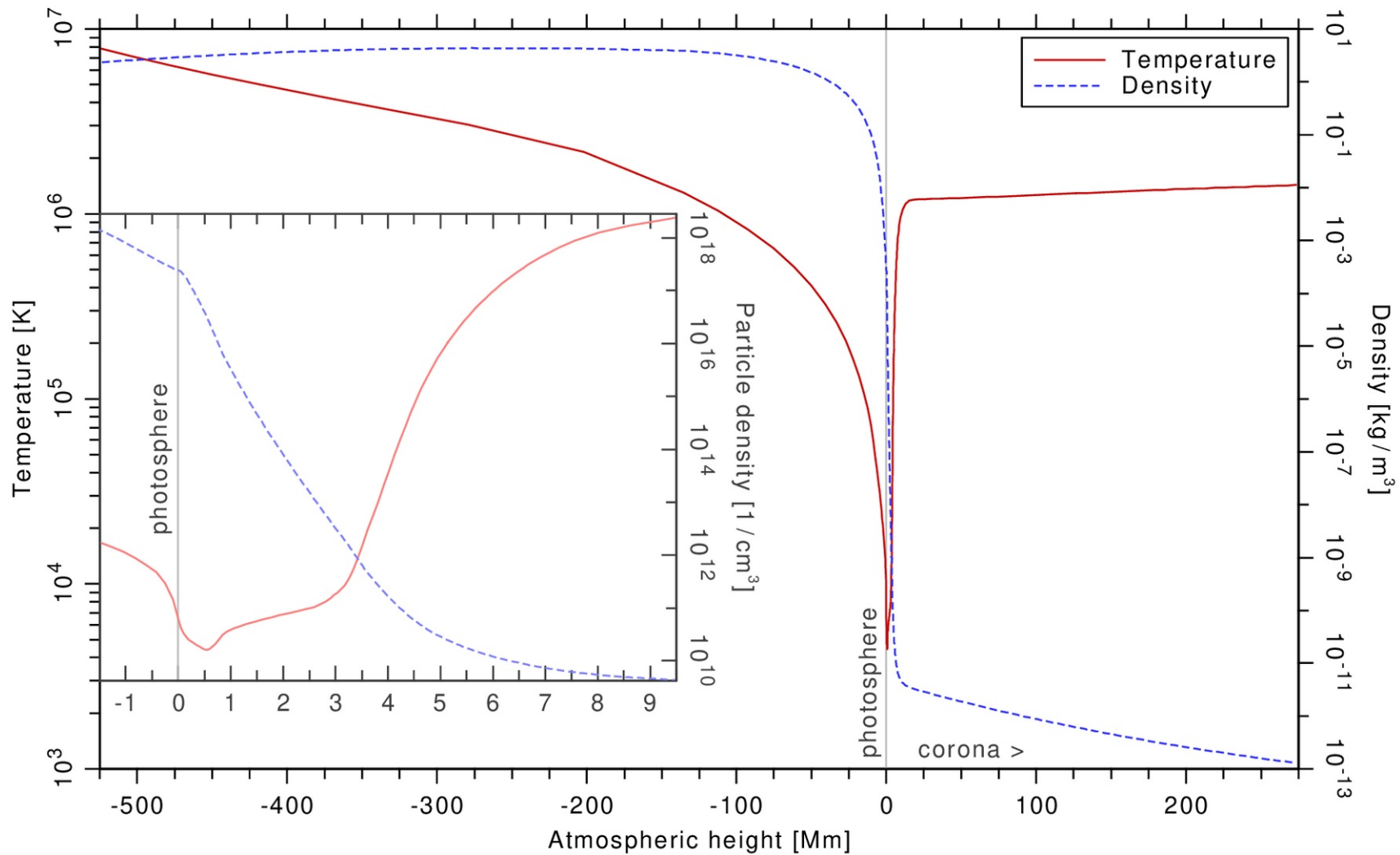
Initial condition

Initial condition

Problem: initial condition might not be in „numerical equilibrium“

Initial condition

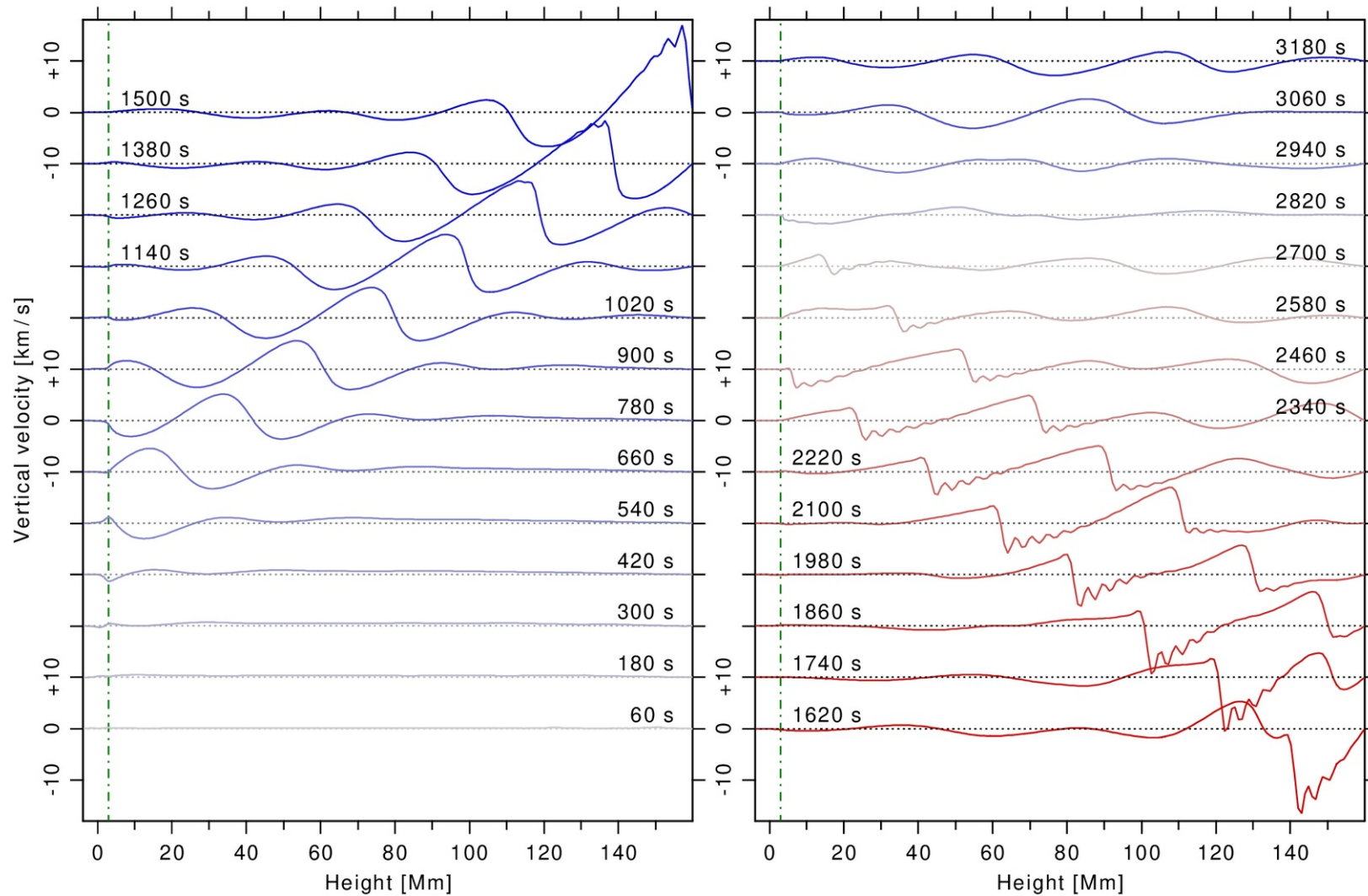
Problem: initial condition might not be in „numerical equilibrium“



=> Solar atmosphere in hydrostatic equilibrium (Bourdin, 2014, CEAB)

Initial condition

Problem: initial condition might not be in „numerical equilibrium“



=> Analytic solution needs time to equilibrate

(Bourdin, 2014, CEAB)

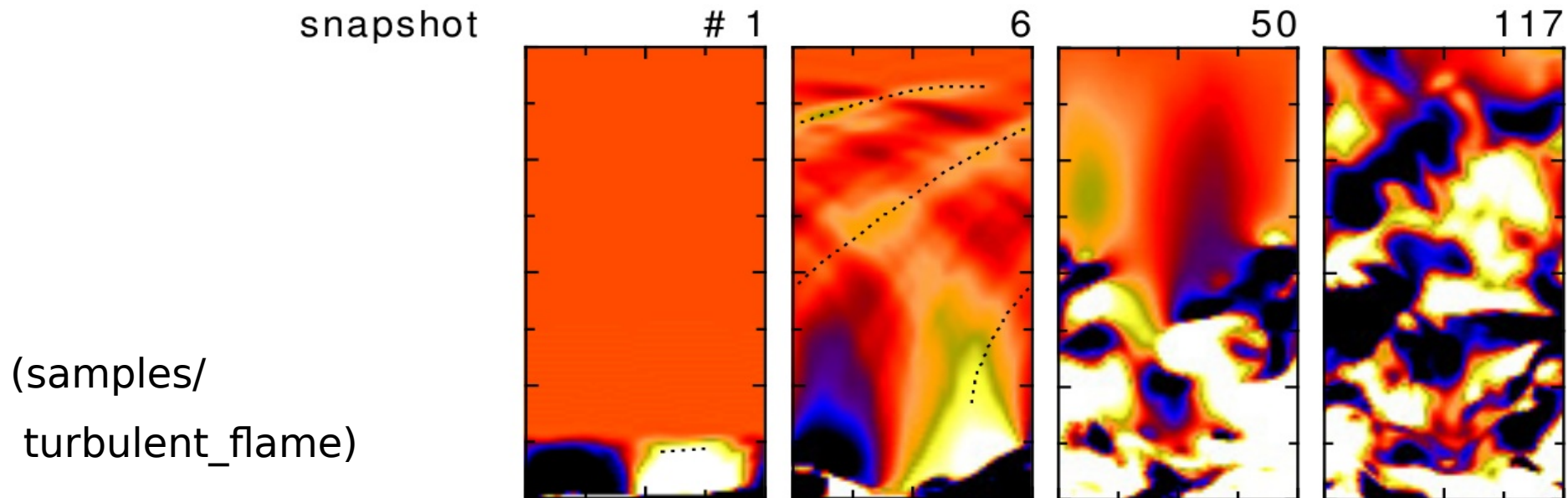
Driving a simulation with external forces

Driving a simulation with external forces

Problem: switching on over small $dt \Rightarrow$ infinite momentum

Driving a simulation with external forces

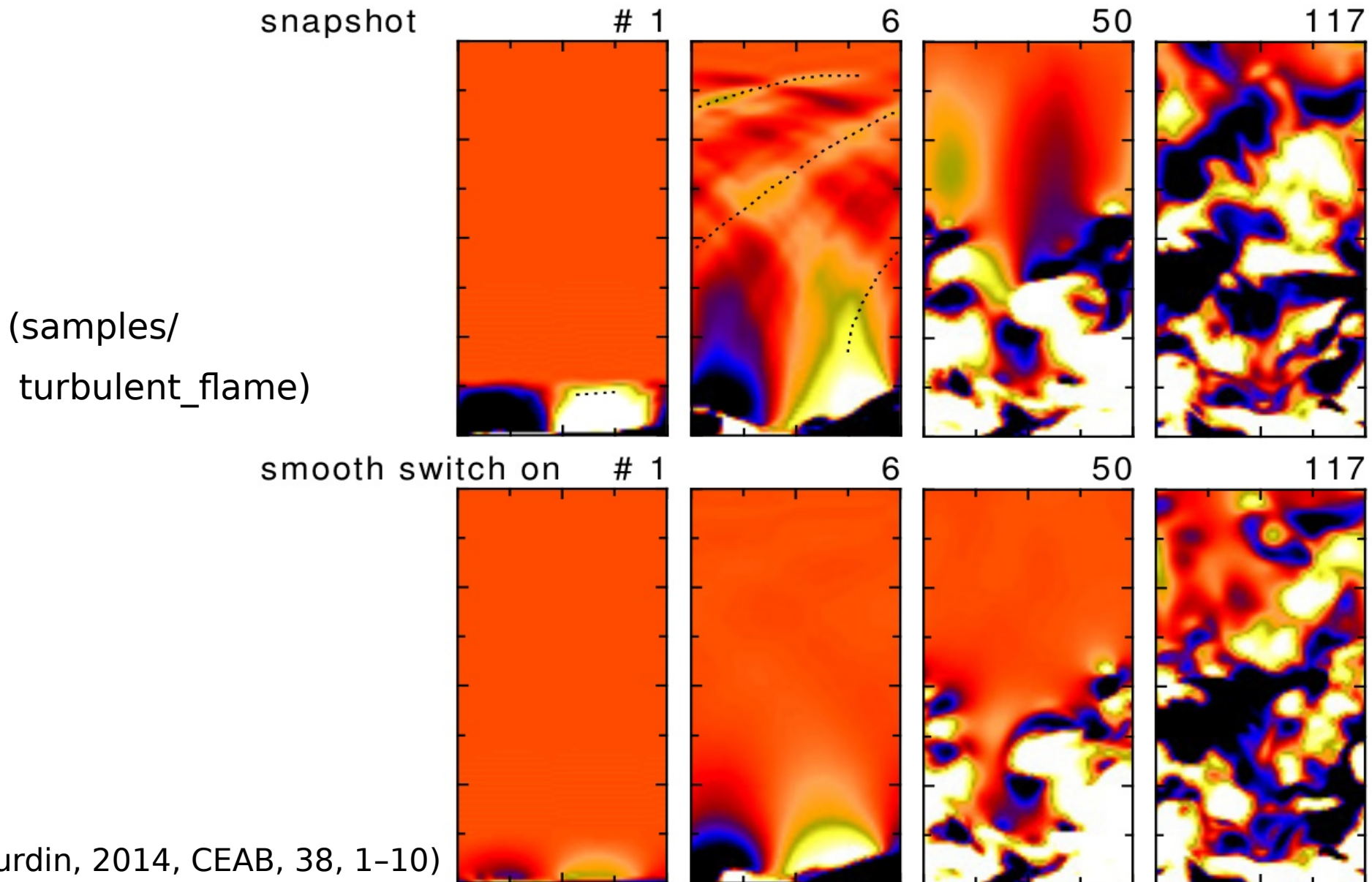
Problem: switching on over small $\Delta t \Rightarrow$ infinite momentum \Rightarrow shock!



(Bourdin, 2014, CEAB, 38, 1-10)

Driving a simulation with external forces

Problem: switching on over small $\Delta t \Rightarrow$ infinite momentum \Rightarrow shock!



(Bourdin, 2014, CEAB, 38, 1-10)